

# **Solver Challenges in Adaptive Mesh Refinement Ice Sheet Modeling**

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## *Joint work with:*

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- ❑ Xylar Asay-Davis (Potsdam-PIK)
- ❑ Stephen Cornford (Bristol)
- ❑ Stephen Price (LANL)
- ❑ Doug Ranken (LANL)
- ❑ Esmond Ng (LBNL)
- ❑ William Collins (LBNL)



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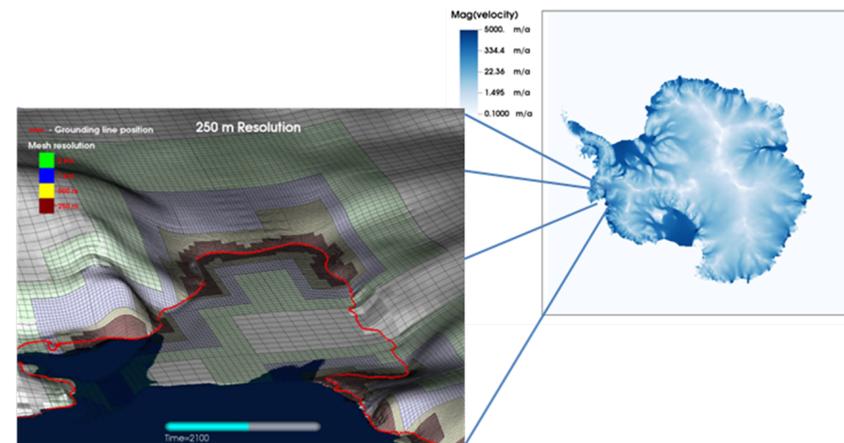
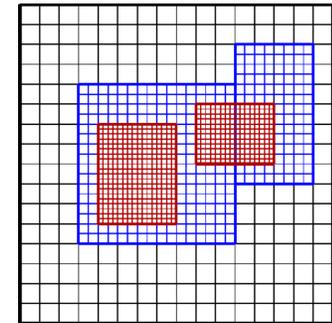
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# BISICLES Ice Sheet Model

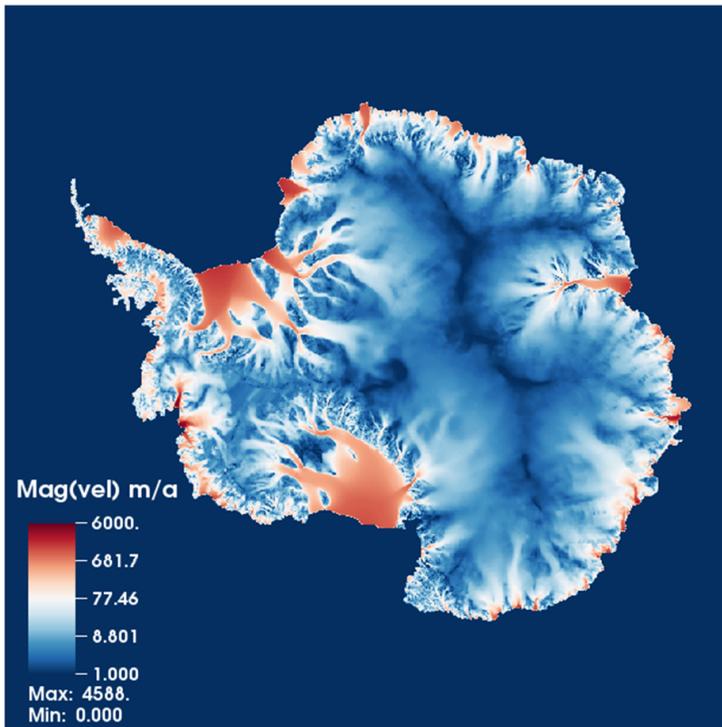
- ❑ Scalable adaptive mesh refinement (AMR) ice sheet model
  - Dynamic local refinement of mesh to improve accuracy
- ❑ Chombo AMR framework for block-structured AMR
  - Support for AMR discretizations
  - Scalable solvers
  - Developed at LBNL
  - DOE ASCR supported (FASTMath)
- ❑ Collaboration with Bristol (U.K.) and LANL
- ❑ Variant of “L1L2” model (Schoof and Hindmarsh, 2010)
- ❑ Coupled to Community Ice Sheet Model (CISM).
- ❑ Users in Berkeley, Bristol, Beijing, Brussels, and Berlin...



# Why is this useful? (another BISICLE for another fish?)



- Ice sheets -- Localized regions where high resolution needed to accurately resolve ice-sheet dynamics (500 m or better at grounding lines)
- Antarctica is really big - too big to resolve at that level of resolution.
- Large regions where such fine resolution is unnecessary (e.g. East Antarctica)
- Well-suited for adaptive mesh refinement (AMR)
- Problems still large: need good parallel efficiency
- Dominated by nonlinear coupled elliptic system for ice velocity solve: good linear and nonlinear solvers



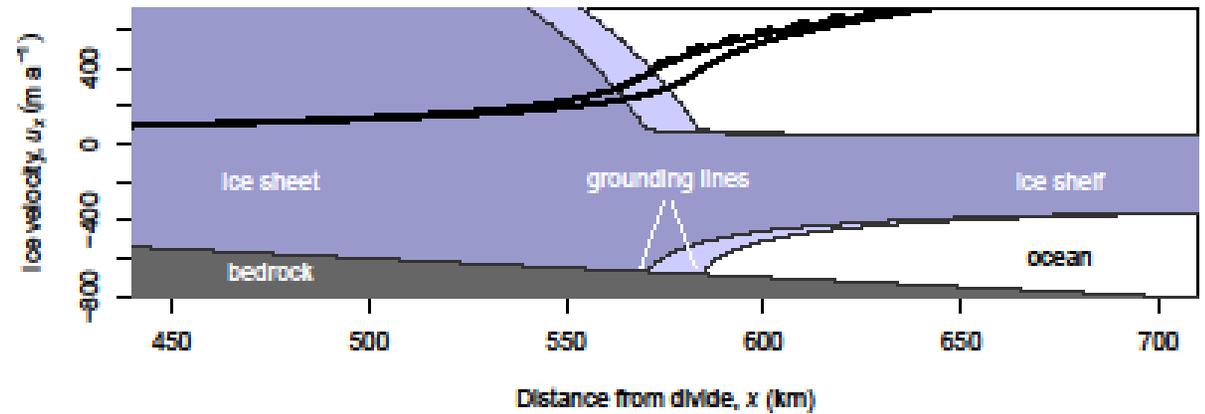
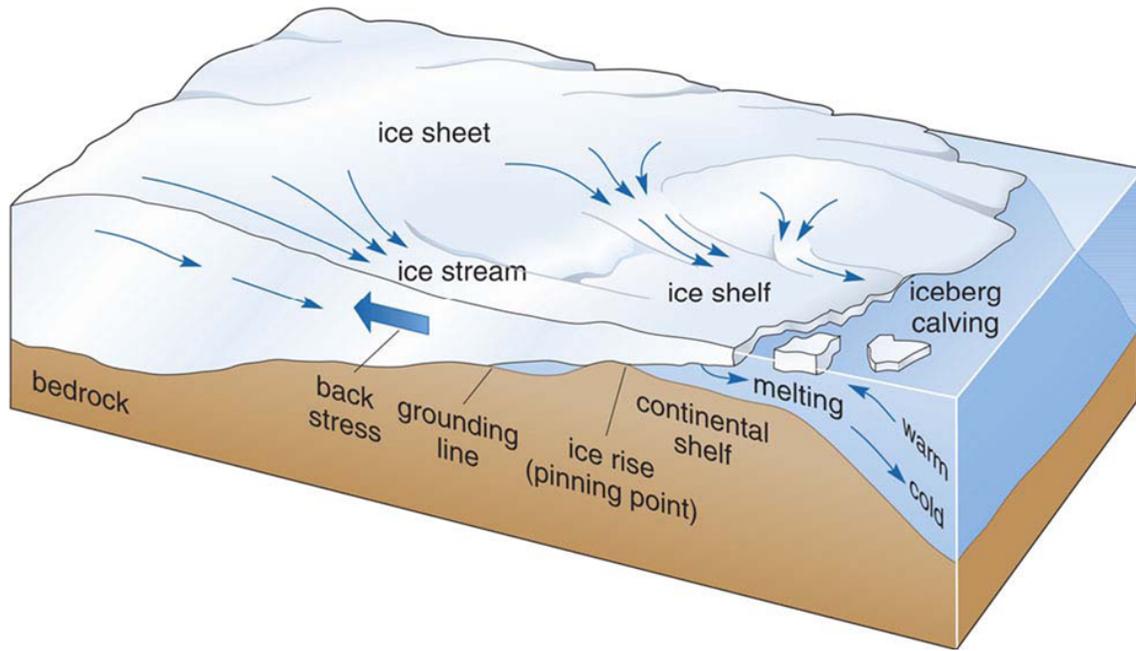
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# Marine Ice Sheets



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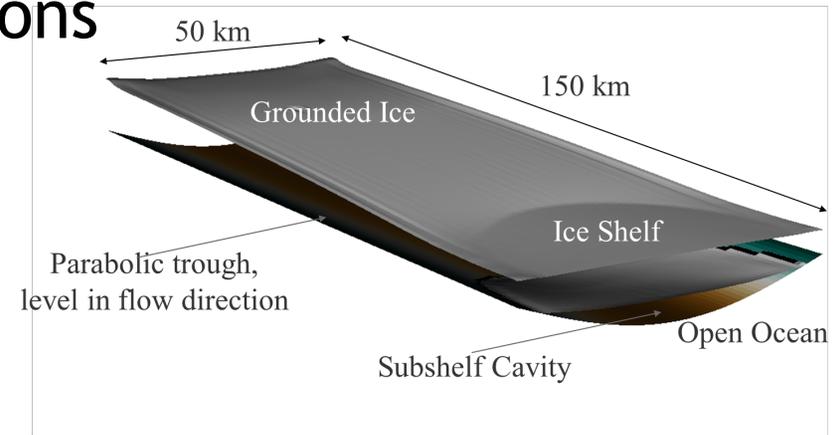
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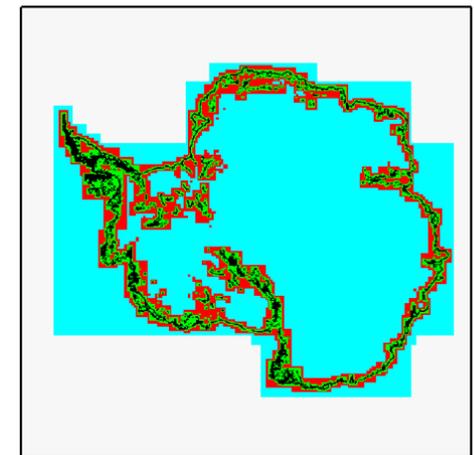
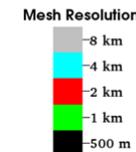
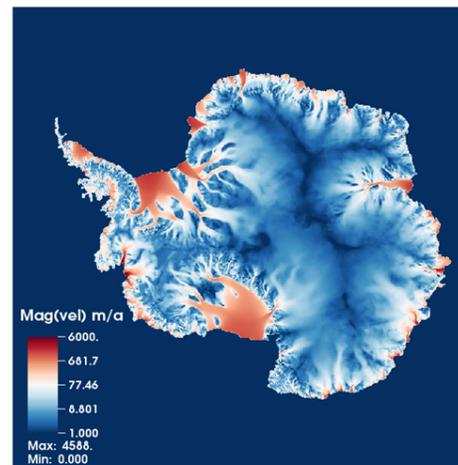
# Target Problems

- ❑ Idealized Ice-Ocean interaction test problems
  - Simple/small geometries designed to understand GL dynamics and ice-ocean interactions

- MISMIP3D
- MISMIP+
- MISOMIP



- ❑ Realistic full-scale
  - Fully-resolved (500m) full-continent
  - Antarctica



# BISICLES: Models and Approximations

**Physics:** Non-Newtonian viscous flow:  $\mu(\dot{\epsilon}, T) = A(T)(\dot{\epsilon})^{(1-n)/2}$

Where  $\dot{\epsilon}$  is the strain rate invariant, typically  $n=3$

## □ Full-Stokes

- Best fidelity to ice sheet dynamics
- Computationally expensive (full 3D coupled nonlinear elliptic equations)

## □ Approximate Stokes

- Use scaling arguments to produce simpler set of equations
- Common expansion is in ratio of vertical to horizontal length scales ( $\epsilon = [h]/[L]$ )
- E.g. Blatter-Pattyn (most common “higher-order” model), accurate to  $O(\epsilon^2)$
- Still 3D, but solve simplified elliptic system (e.g. 2 coupled equations)

## □ Depth-integrated

- “Shallow Ice” and “Shallow-Shelf” approximations (accurate to  $O(\epsilon)$ )

▪ Special case of approximate Stokes with 2D equation set



# “L1L2” Model (Schoof and Hindmarsh, 2010).

- Uses asymptotic structure of full Stokes system to construct a higher-order approximation
  - Expansion in  $\epsilon = [H]/[L]$  and  $\lambda = [\tau_{\downarrow shear}]/[\tau_{\downarrow normal}]$  (ratio of shear & normal stresses)
    - Large  $\lambda$ : shear-dominated flow
    - Small  $\lambda$ : sliding-dominated flow
  - Computing velocity to  $O(\epsilon^2)$  only requires  $\tau$  to  $O(\epsilon)$
- Computationally **much** less expensive -- enables fully 2D vertically integrated discretizations. (can reconstruct 3d)
- Similar formal accuracy to Blatter-Pattyn  $O(\epsilon^2)$ 
  - Recovers proper fast- and slow-sliding limits:
    - SIA ( $1 \ll \lambda \leq \epsilon^{-1/n}$ ) -- accurate to  $O(\epsilon^2 \lambda^{n-2})$
    - SSA ( $\epsilon \leq \lambda \leq 1$ ) - accurate to  $O(\epsilon^2)$

# “L1L2” Model (Schoof and Hindmarsh, 2010), cont.

- Can construct a computationally efficient scheme:
  1. Approximate constitutive relation relating  $grad(u)$  and stress field  $\tau$  with one relating  $grad(u|_{z=b})$ , vertical shear stresses  $\tau_{xz}$  and  $\tau_{yz}$  given by the SIA / lubrication approximation and other components  $\tau_{xx}(x,y,z)$ ,  $\tau_{xy}(x,y,z)$ , etc
  2. leads to an effective viscosity  $\mu(x,y,z)$  which depends only on  $grad(u|_{z=b})$  and  $grad(z|_s)$ , ice thickness, etc
  3. Momentum equation can then be integrated vertically, giving a nonlinear, 2D, elliptic equation for  $u|_{z=b}(x,y)$
  4.  $u(x,y,z)$  can be reconstructed from  $u|_{z=b}(x,y)$



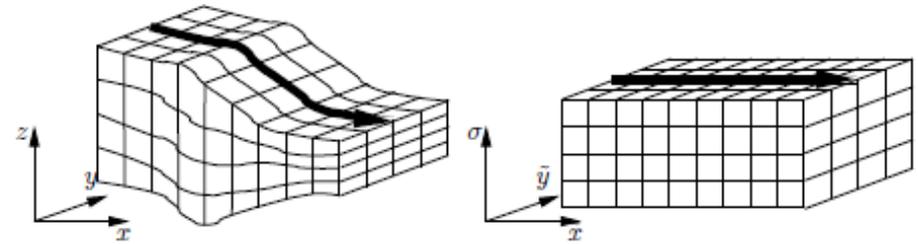
# Modified “L1L2” Model (SSA\*)

- Can construct a computationally efficient scheme:
  1. Approximate constitutive relation relating  $grad(u)$  and stress field  $\tau$  with one relating  $grad(u|_{z=b})$ , vertical shear stresses  $\tau_{xz}$  and  $\tau_{yz}$  given by the SIA / lubrication approximation and other components  $\tau_{xx}(x,y,z)$ ,  $\tau_{xy}(x,y,z)$ , etc
  2. leads to an effective viscosity  $\mu(x,y,z)$  which depends only on  $grad(u|_{z=b})$  and  $grad(z|_s)$ , ice thickness, etc
  3. Momentum equation can then be integrated vertically, giving a nonlinear, 2D, elliptic equation for  $u|_{z=b}(x,y)$
  4.  ~~$u(x,y,z)$  can be determined from  $u|_{z=b}(x,y)$~~
  4. Use  $u(x,y,z) = u|_{z=b}(x,y)$  (neglect vertical shear in flux velocity)



# Discretizations

- Baseline model:
  - Logically-rectangular grid, obtained from a time-dependent uniform mapping.
  - 2D equation for ice thickness



$$\frac{\partial H}{\partial t} = b - \nabla \cdot H\bar{u}$$

- Vertically-integrated momentum balance results in 2D nonlinear viscous tensor solve (viscosity a function of velocity) for velocity  $u \downarrow b$  at the base of the ice:

$$\beta \nabla \cdot u \downarrow b + \nabla \cdot [\mu(\varepsilon \nabla \cdot u \downarrow b)(\nabla + \nabla \nabla \cdot T)u \downarrow b - 2\mu(\nabla \cdot u \downarrow b)] = -g/\rho H \nabla s$$

$\beta \nabla \cdot u \downarrow b$  = friction coefficient,  $\varepsilon$  = strain rate invariant of ice velocity,  $g$  = gravity,  $\rho$  = ice density,  $H$  = ice thickness,  $\nabla s$  = horizontal gradient of upper surface

- Enthalpy formulation (Aschwanden) for energy



# Discretizations, cont

- ❑ Use of Finite-volume discretizations (vs. Finite-difference discretizations) simplifies implementation of local refinement.
- ❑ Software implementation based on constructing and extending existing solvers using the Chombo libraries.



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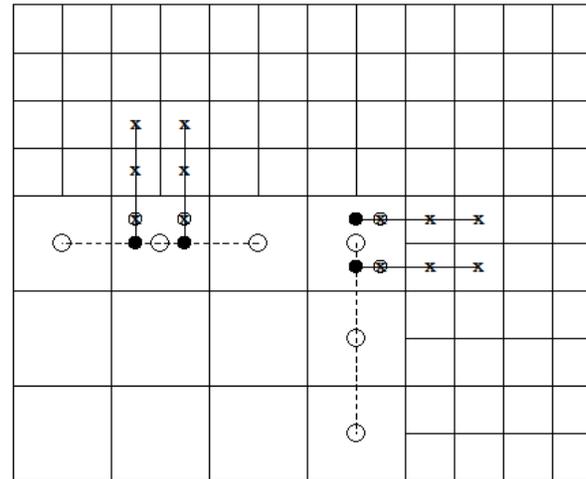
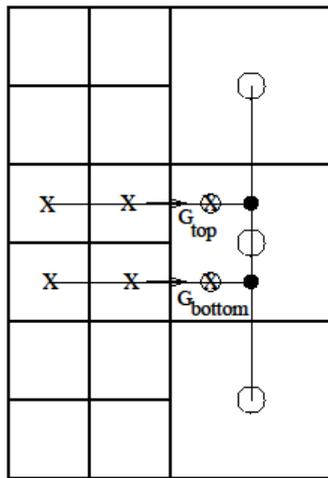
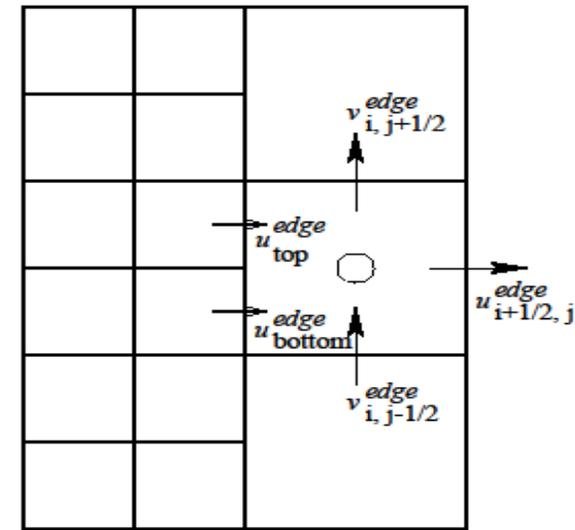


# AMR and Finite-Volumes

- AMR and finite-volume discretizations play well together:

$$L(\varphi) = \nabla \cdot F$$

- Complex discretizations at coarse-fine interfaces:



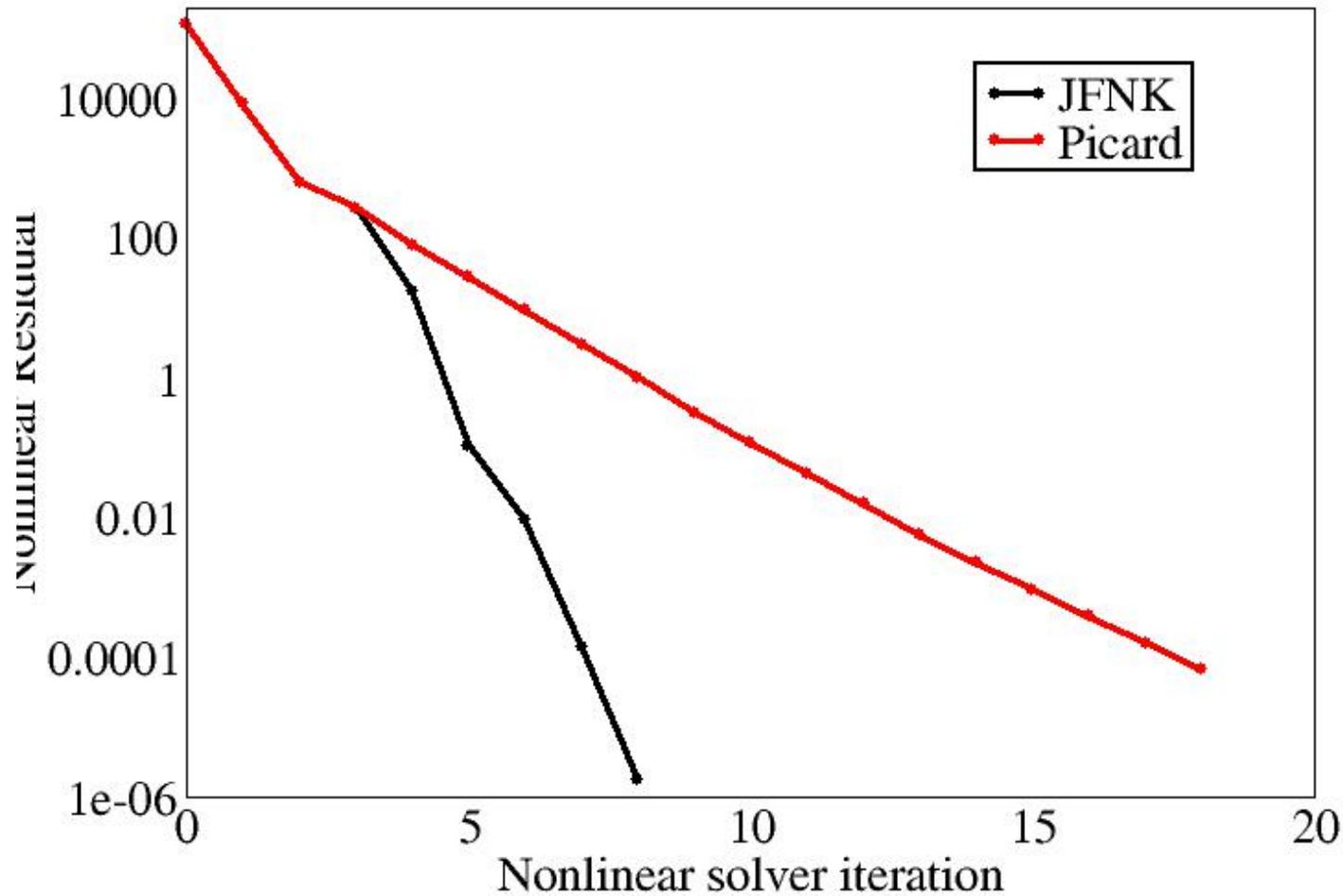
# Nonlinear Solvers

- ❑ Most computational effort spent in nonlinear ice velocity solve.
- ❑ Picard iteration:
  - Robust
  - Simple to implement
  - Slow (but steady) convergence
- ❑ Jacobian-free Newton-Krylov (JFNK):
  - More complex to implement
  - Works best with decent initial guess
  - Rapid convergence
  - Well-suited for Chombo AMR elliptic solvers
- ❑ Approach - Picard iteration initially, switch to JFNK when convergence slows



# Nonlinear Solvers (cont)

Nonlinear Solver Convergence



# Initialization

- Glen's Law singularity ( $\frac{1}{\epsilon^{12}}$ ) as  $\epsilon \uparrow \rightarrow 0$ 
  - Regularization: use  $(\epsilon^{12} + \delta)$  in Glen's law ( $\delta = O(10^{-12})$ )
- Natural initial guess ( $u=0$ ) too close to singularity
  - Takes a lot of solver effort/iterations to push solution away from singularity
  - Better idea - initial linear (constant  $\frac{1}{\epsilon^{12}}$ ) solve
    - Relatively inexpensive
    - Pushes solution away from singularity -- reasonable initial guess for nonlinear solve
    - (equivalent to homotopy/regularization with a large  $\frac{1}{\epsilon^{12}}$ )
- Natural form of grid sequencing from AMR



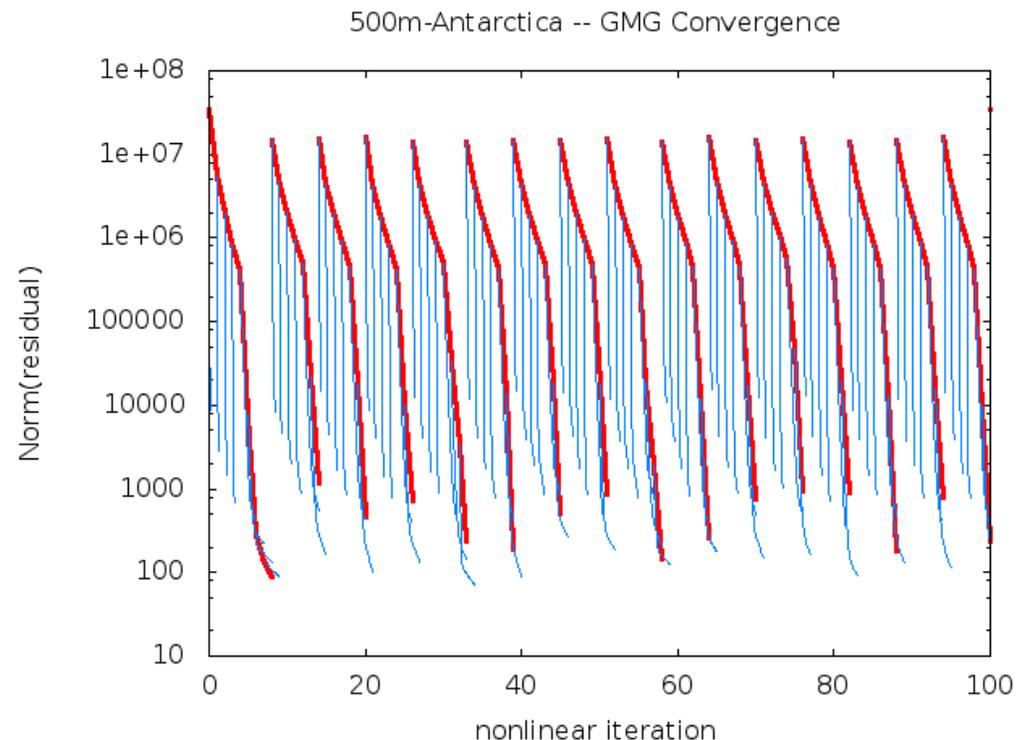
# Well-posedness

- ❑ Disconnected floating ice (icebergs) are ill-posed
- ❑ Can emerge through Ice dynamics (ice-ocean coupling)
- ❑ Can also emerge from AMR regridding - interpolation
- ❑ Solution - sweep domain for disconnected ice
  - “Marching” scheme (start with grounded ice and march outward in repeated passes)
  - AMR-aware connected-components scheme (scalable)
    - (Zou, Martin, et al (2015) )



# Linear Solvers

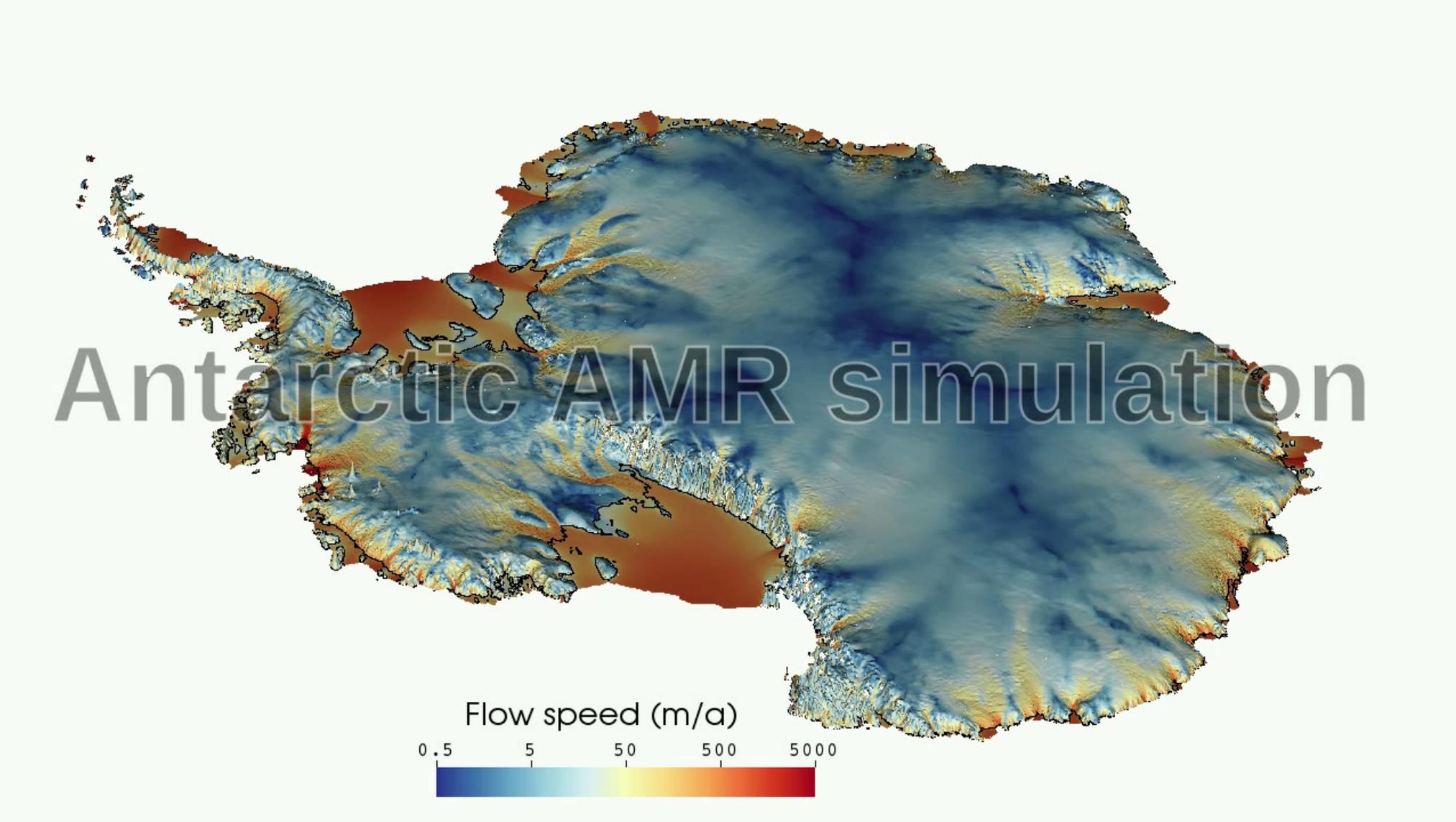
- ❑ Need good linear solver performance!
- ❑ Chombo native solvers - Geometric MultiGrid (GMG)
  - Follows Naturally from AMR hierarchy
- ❑ When it works, it works really well (after some tuning)
  - Matrix-free!
  - Relatively efficient
- ❑ Works well enough on some problems (Full-continent Antarctica!)



# Example - 1000-year Antarctic simulations

- ❑ Range of finest resolution from 8 km (no refinement) to 500m (4 levels of factor-2 refinement)
- ❑ At initial time, subject ice shelves to extreme (outlandish) melting:
  - No melt for  $h < 100\text{m}$
  - Range up to 800m/a where  $h > 400\text{m}$ .
  - **No melt applied in partially-grounded cells**
- ❑ For each resolution, evolve for 1000 years
- ❑ Solver tolerances relatively loose

# Results:



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# Linear Solvers, cont.

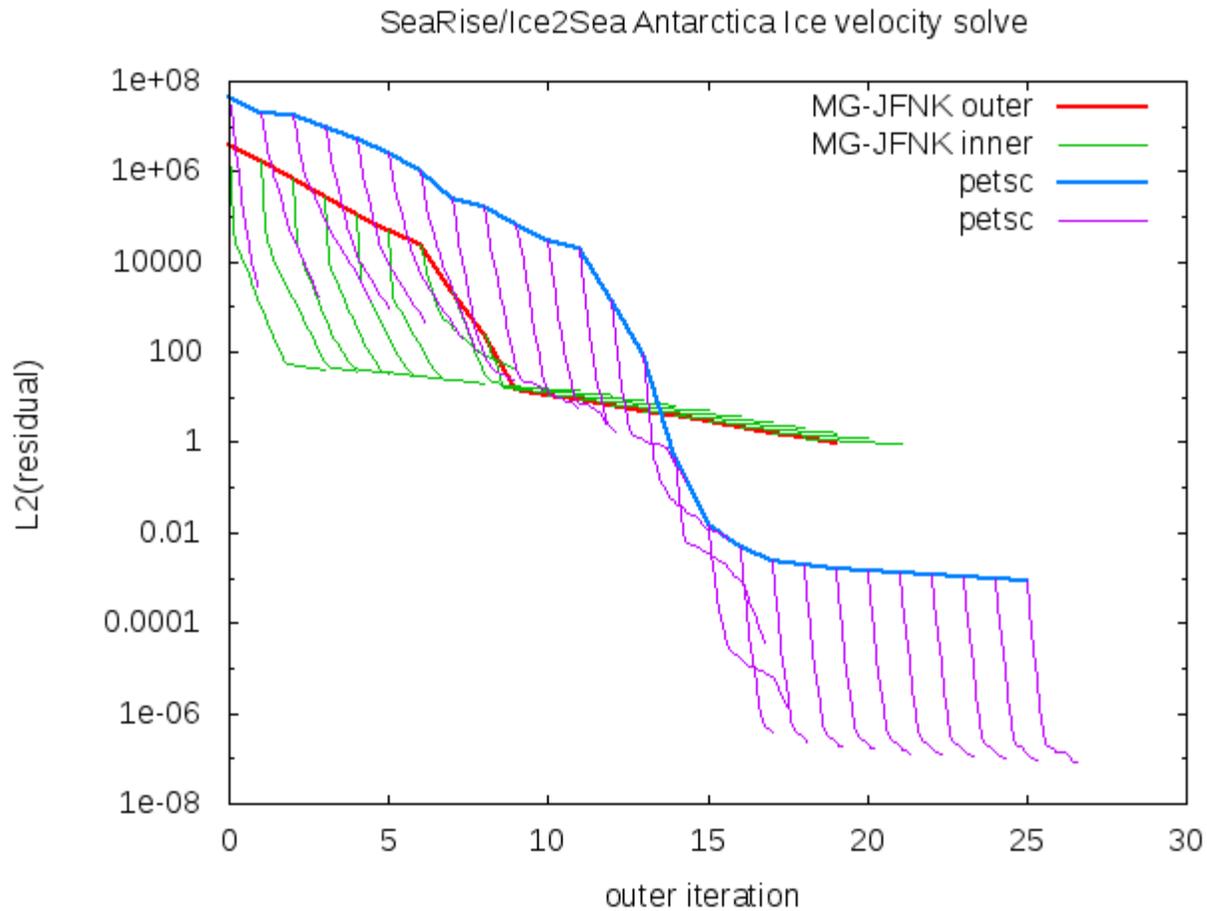
- ❑ **However** - GMG sometimes fails (stalls, diverges)
  - Stalling - finer than 1 km resolution
  - Failure - where Ice shelves are important (Ice/ocean)
- ❑ Sharp/strong coefficient gradients (hard to coarsen)
- ❑  $\beta \ll 1$  (Change from Parabolic to Elliptic) at GL:

$$\beta \nabla^2 u + \nabla \cdot [\mu(\epsilon \nabla^2) (\nabla + \nabla \cdot T) u] - 2\mu (\nabla \cdot u) = -g/\rho \quad H \nabla s$$

- ❑ Solution - AMG (PETSc gamg/Hypre BoomerAMG)



# Linear Solvers - GAMG vs. Geometric MG



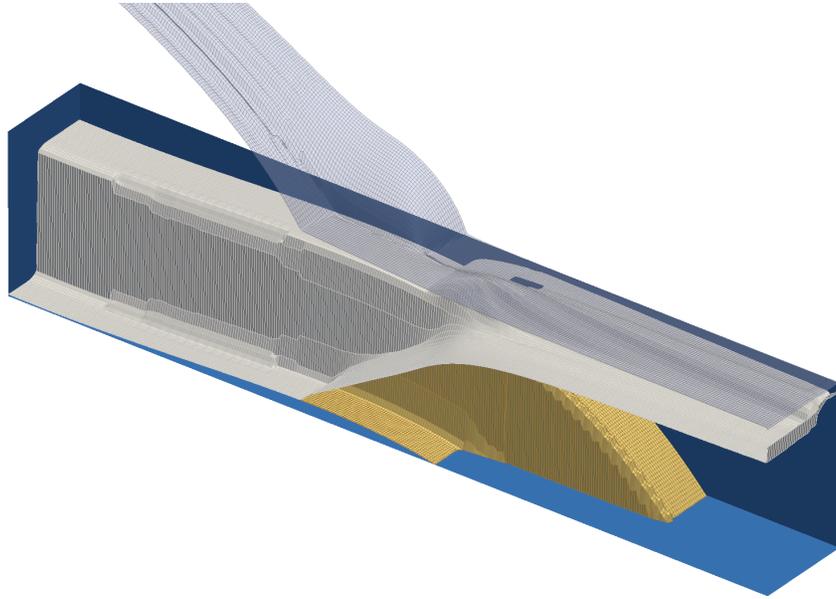
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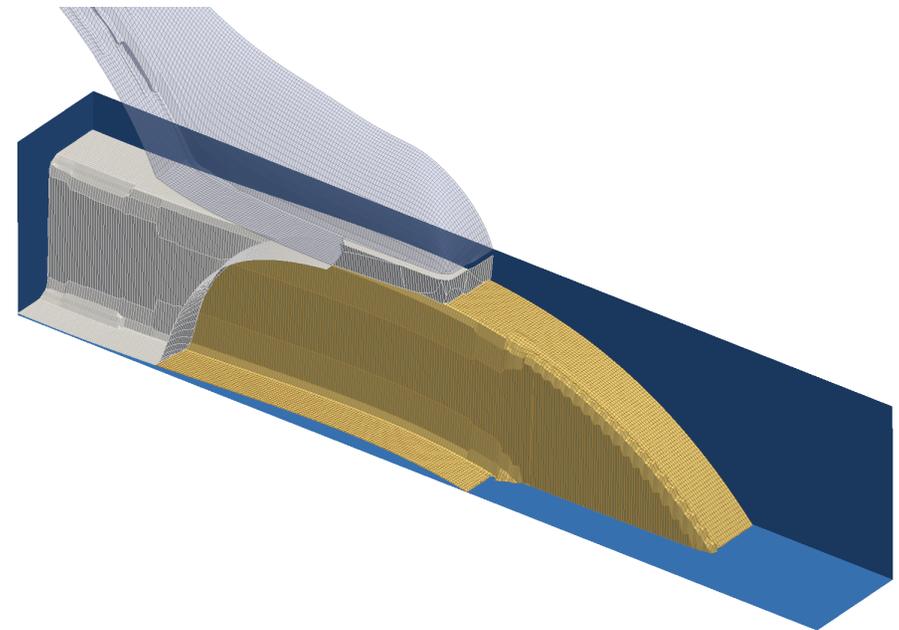
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# MISOMIP (Asay-Davis et al (2015) )



Steady-state initial condition



Fully-retreated condition

- Marine Ice Sheet-Ocean Model Intercomparison Project
- Ice Sheet coupled to Ocean Model through melt rates
- Driven by far-field forcing -
  - $0 < t < 100$  years: Warm Phase (1 C)
  - $100 < t < 200$  years: Cold Phase (-1.9 C)



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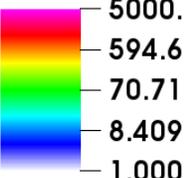
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# MISOMIP

DB: plot.misomip.000000.2d.hdf5  
Cycle: 0      Time: 0.00274

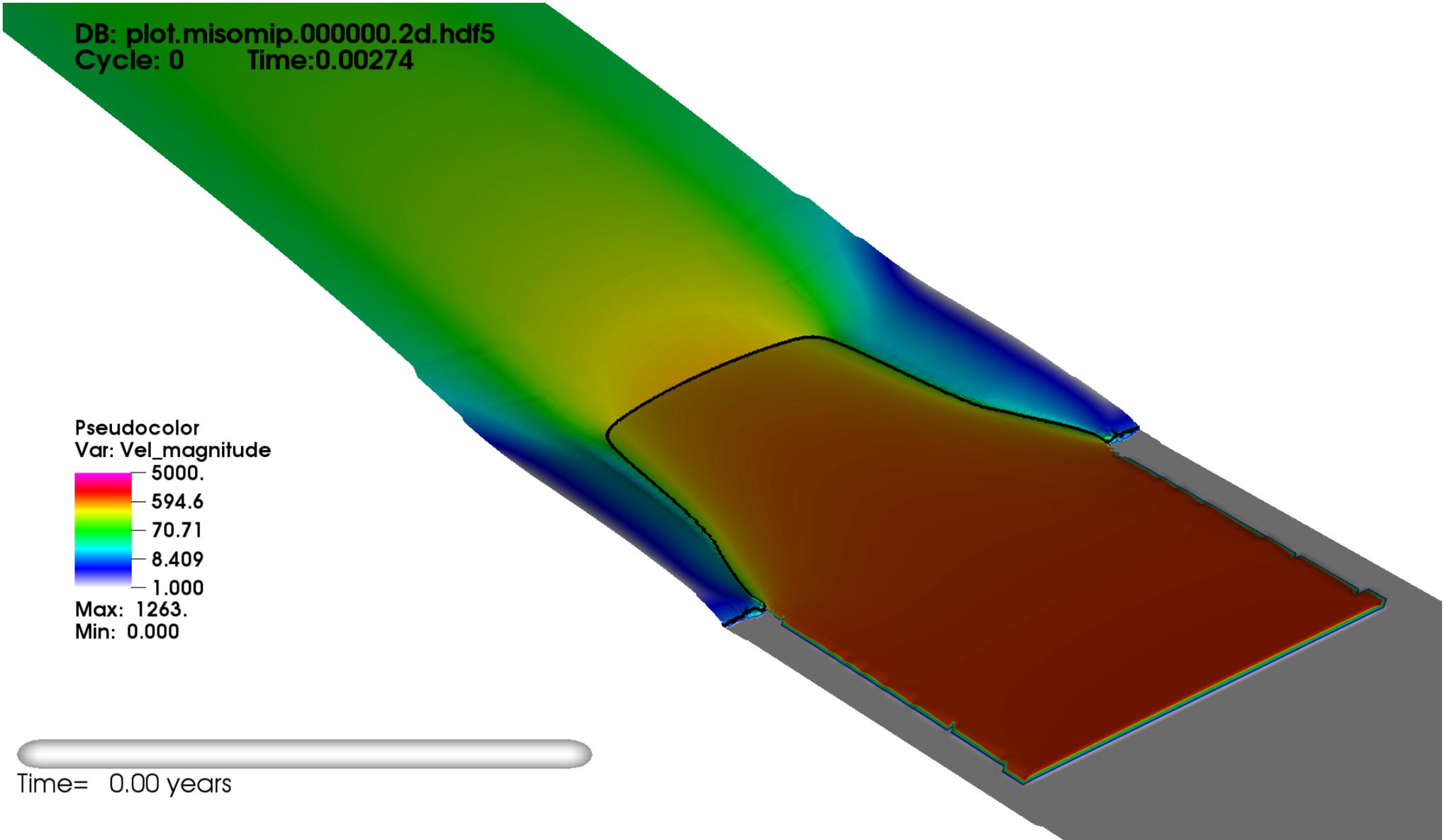
Pseudocolor  
Var: Vel\_magnitude



5000.  
594.6  
70.71  
8.409  
1.000

Max: 1263.  
Min: 0.000

Time= 0.00 years



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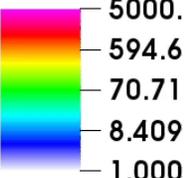
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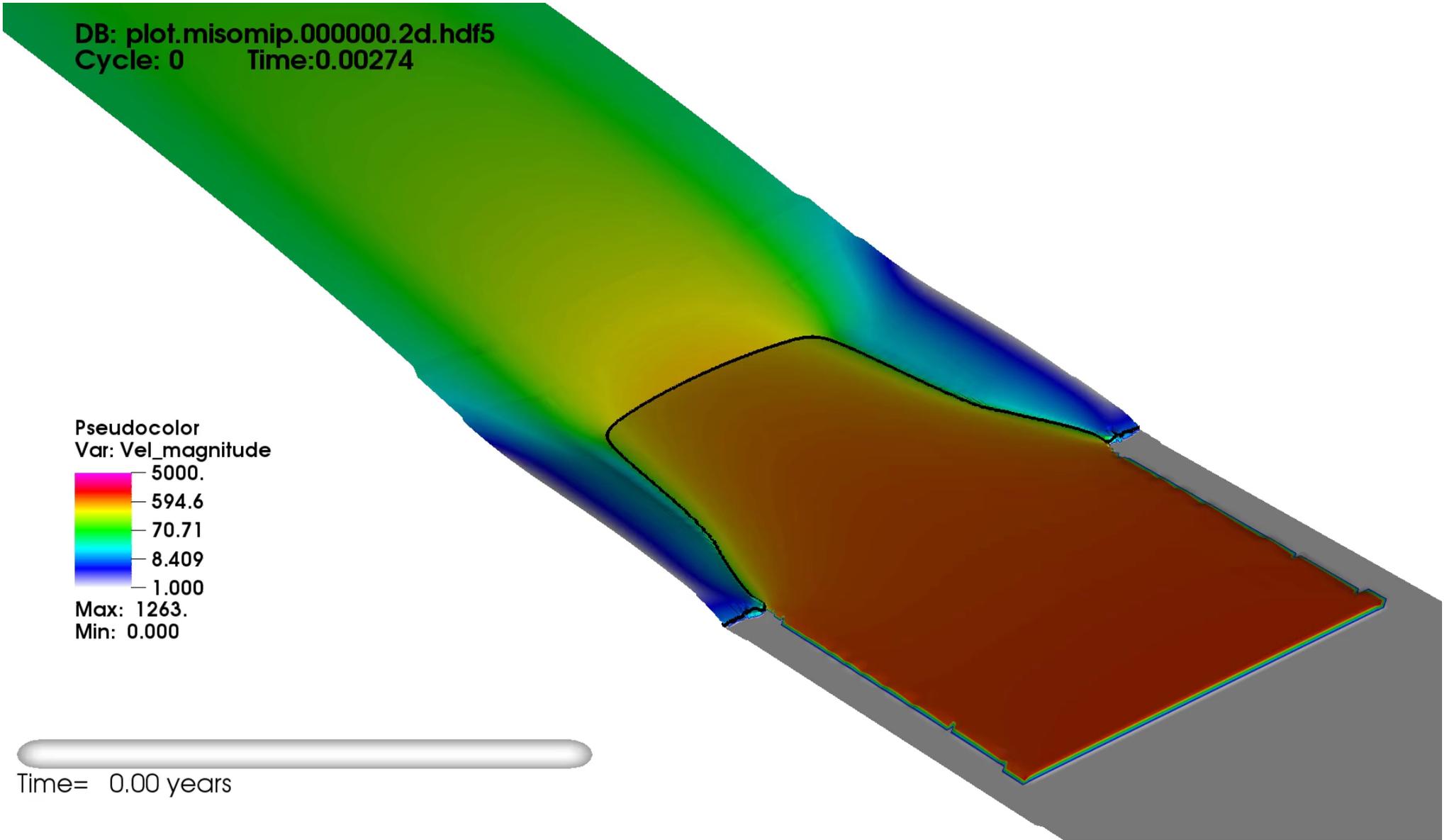
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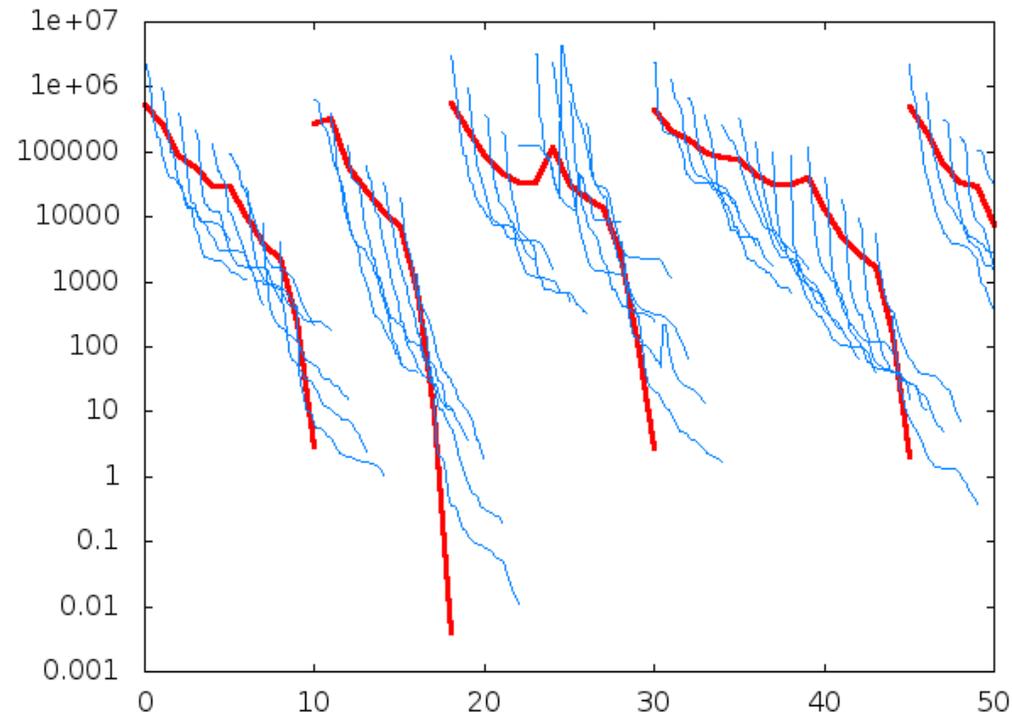
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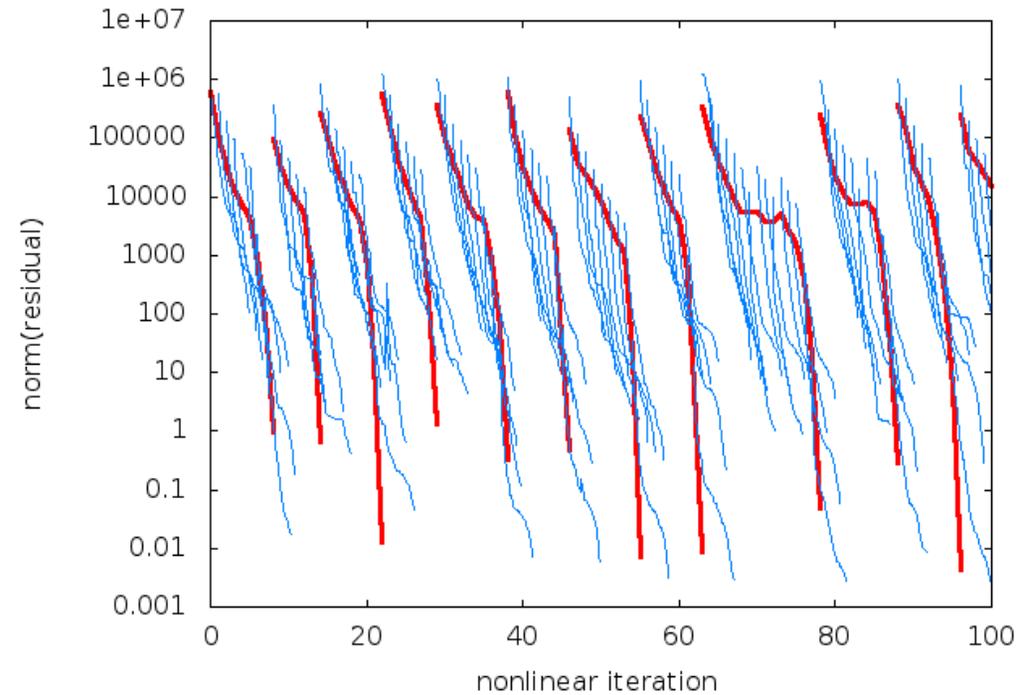


# MISOMIP - Solver Convergence

MISOMIP-IceOcean2 year 50 Solver Convergence



MISOMIP-IceOcean2 year 85 Solver Convergence



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# *FAS Multigrid nonlinear solver*

- ❑ Full Approximation Storage (FAS) - nonlinear multigrid
- ❑ Picard, JFNK:
  - linear solver nested inside of nonlinear one
  - Linear Multigrid solvers (residual-correction form) work well.
- ❑ FAS Multigrid - fully nonlinear solver (no outer solver)
  - Can outperform JFNK/MG
  - More robust (don't need good initial guess)
  - Simpler to implement and maintain
  - Nonlinear convergence similar to MG linear convergence



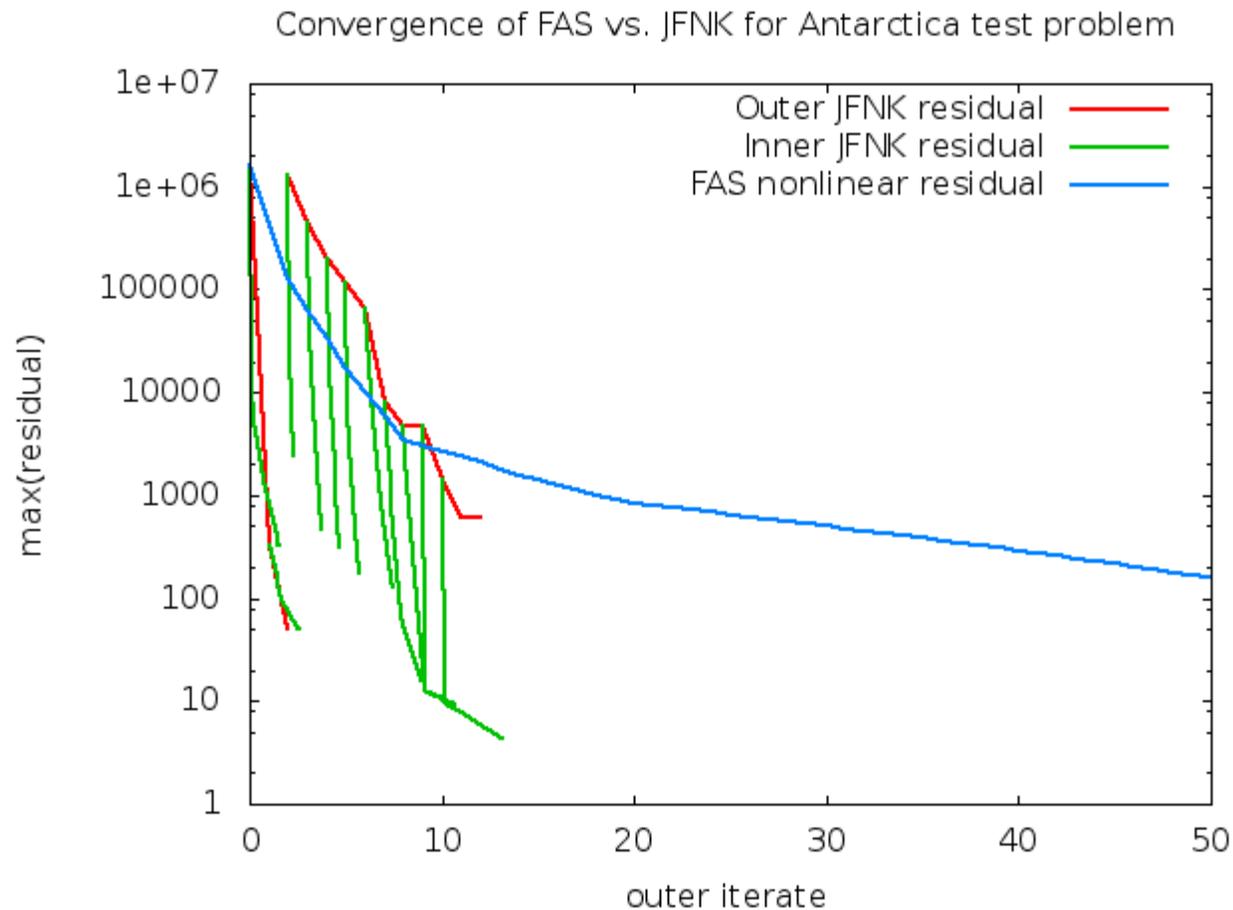
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# FAS Multigrid nonlinear solver (cont)



Solution time for 8 cores on linux desktop:  
FAS-MG: 715 s      JFNK-MG: 1128 s



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- ❑ NERSC



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# Extras

# Grounding-line dynamics experiments

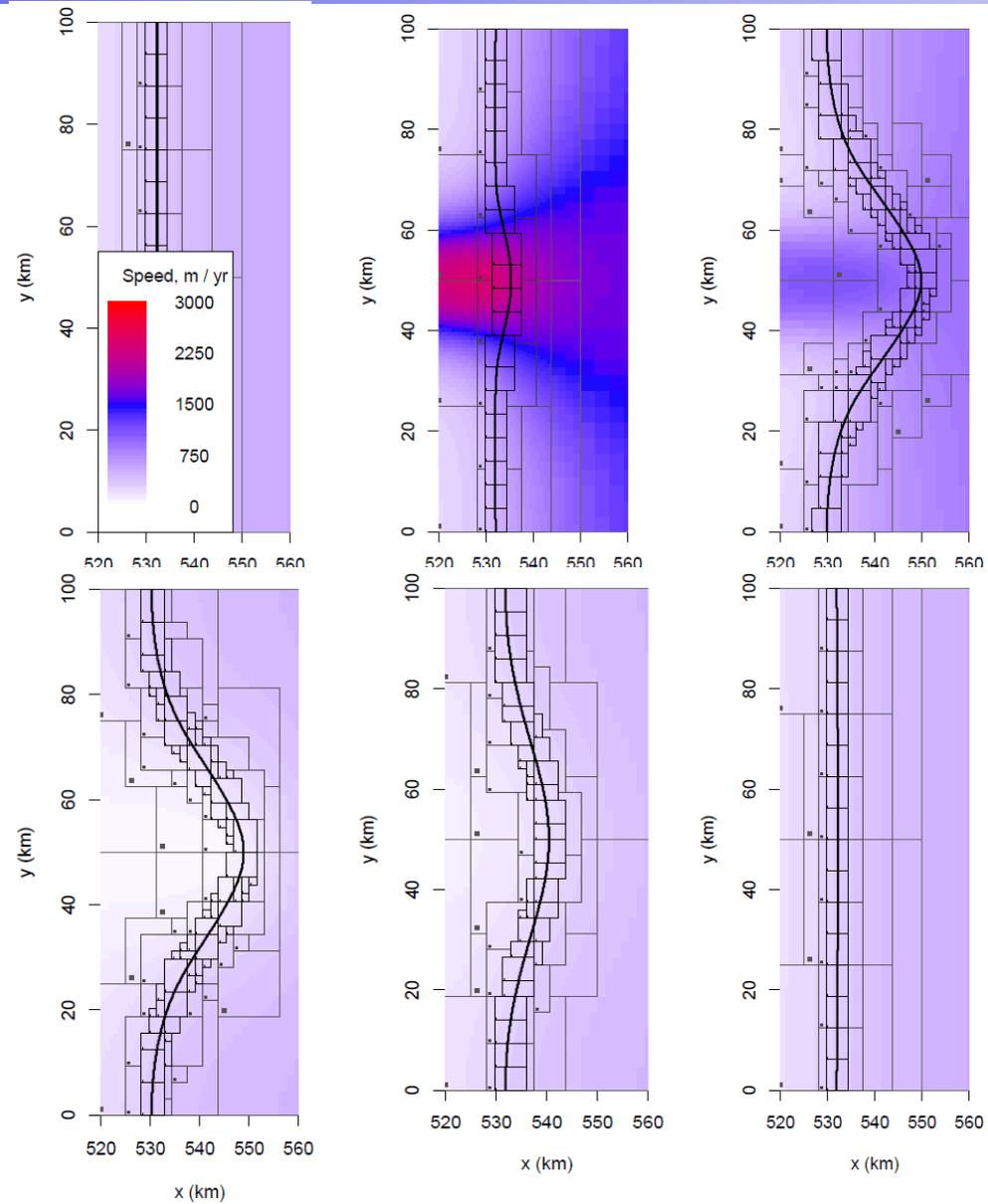
- ❑ Series of ice-sheet modeling community model intercomparison projects designed to understand issues in modeling of GLs
  - MISMIP, MISMIP3D, MISMIP+
- ❑ All point to a need for very fine spatial resolution to get GL dynamics right (sub-km in most cases)
- ❑ Prime use case for adaptive mesh refinement (AMR)



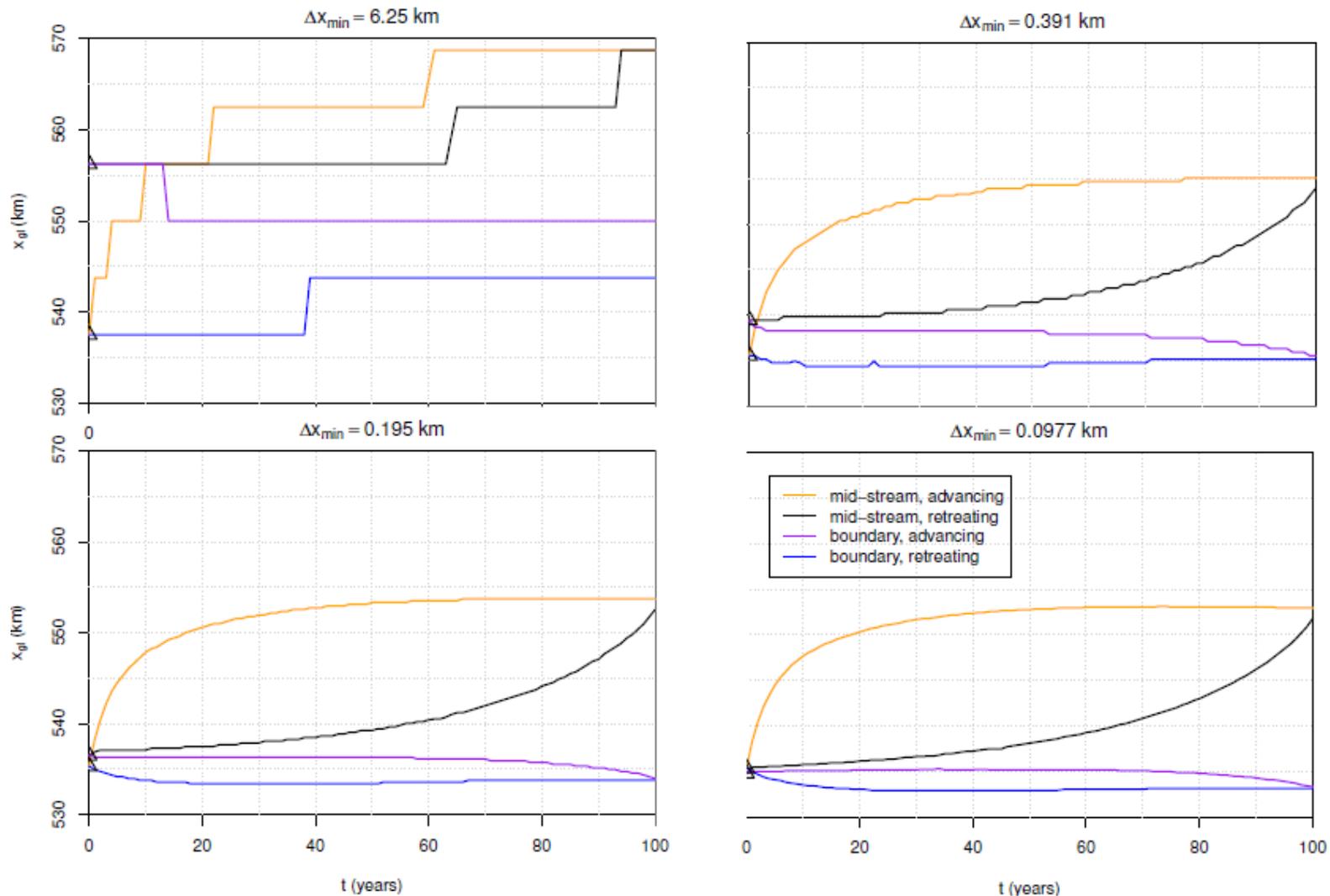
# BISICLES Results - MISMIP3D

## Experiment P75R: (Pattyn et al (2011))

- ❑ Begin with steady-state (equilibrium) grounding line.
- ❑ Add Gaussian slippery spot perturbation at center of grounding line
- ❑ Ice velocity increases, GL advances.
- ❑ After 100 years, remove perturbation.
- ❑ Grounding line should return to original steady state.
- ❑ Figures show AMR calculation:
  - $\Delta x \downarrow 0 = 6.5 \text{ km}$  base mesh,
  - 5 levels of refinement
  - Finest mesh  $\Delta x \downarrow 4 = 0.195 \text{ km}$ .
  - $t = 0, 1, 50, 101, 120, 200 \text{ yr}$
- ❑ Boxes show patches of refined mesh.
- ❑ GL positions match Elmer (full-Stokes)



# MISMIP3D (cont): L1L2 (SSA\*) Spatial Resolution

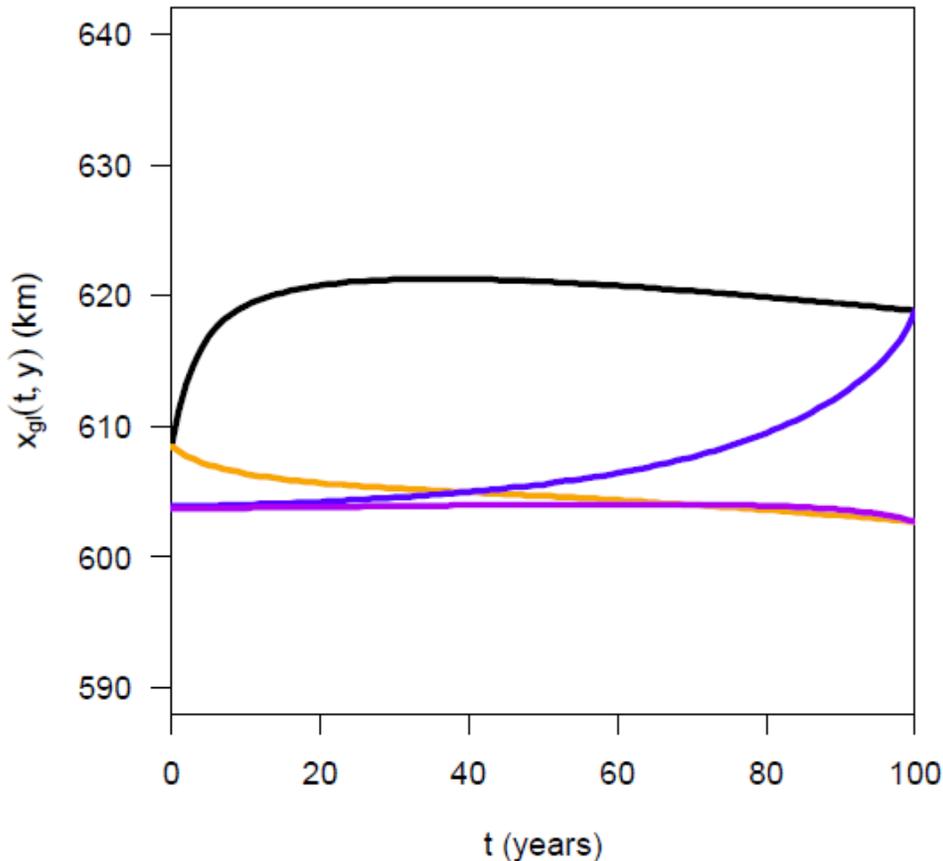


- Very fine (~200 m) resolution needed to achieve reversability!

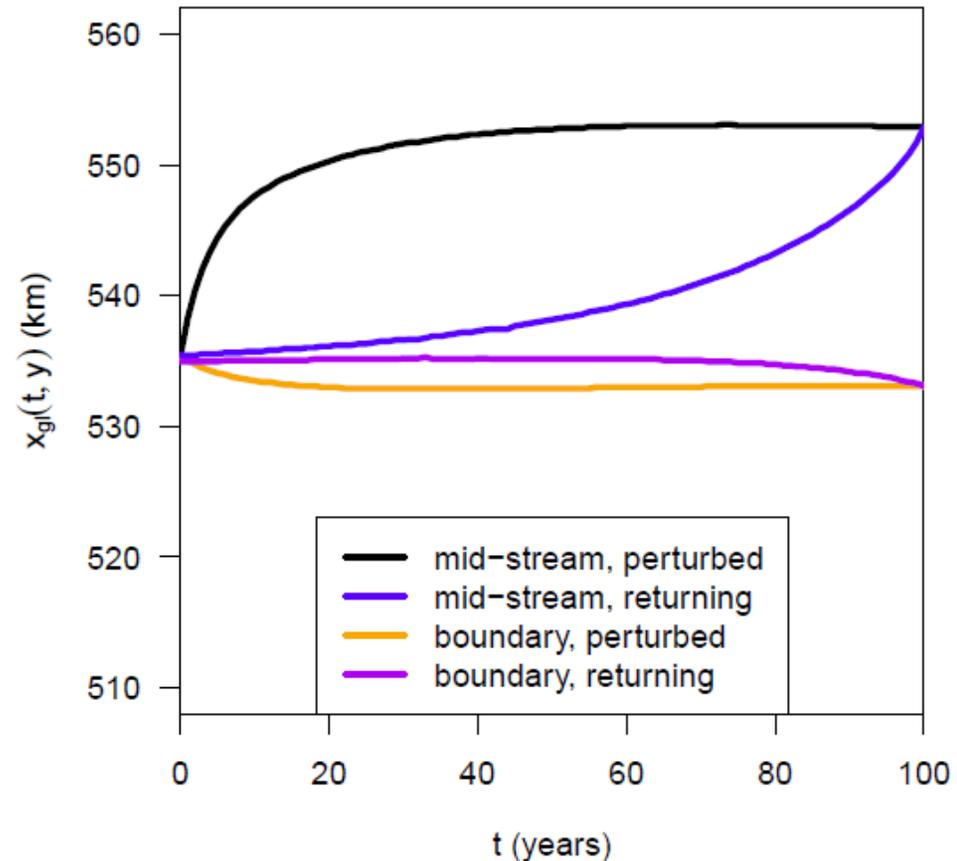


# MISMIP3D: SSA vs. “L1L2” or “SSA\*”

SSA,  $\Delta x^L = 100$  m



SSA\*,  $\Delta x^L = 100$  m



- Direct comparison of SSA vs. SSA\*
  - (fully resolved spatially, same numerics, etc)
  - Note difference in steady-state GL positions



# ***BISICLES Results - Ice2Sea Amundsen Sea***

- ❑ Study of effects of warm-water incursion into Amundsen Sea.
- ❑ Results from Payne et al, (2012), submitted.
- ❑ Modified 1996 BEDMAP geometry (Le Brocq 2010), basal traction and damage coefficients to match Joughin 2010 velocity.
- ❑ Background SMB and basal melt rate chosen for initial equilibrium.
- ❑ SMB held fixed.
- ❑ Perturbations in the form of additional subshelf melting:
  - derived from FESOM circumpolar deep water
  - ~5 m/a in 21<sup>st</sup> Century,
  - ~25 m/a in 22<sup>nd</sup> Century.



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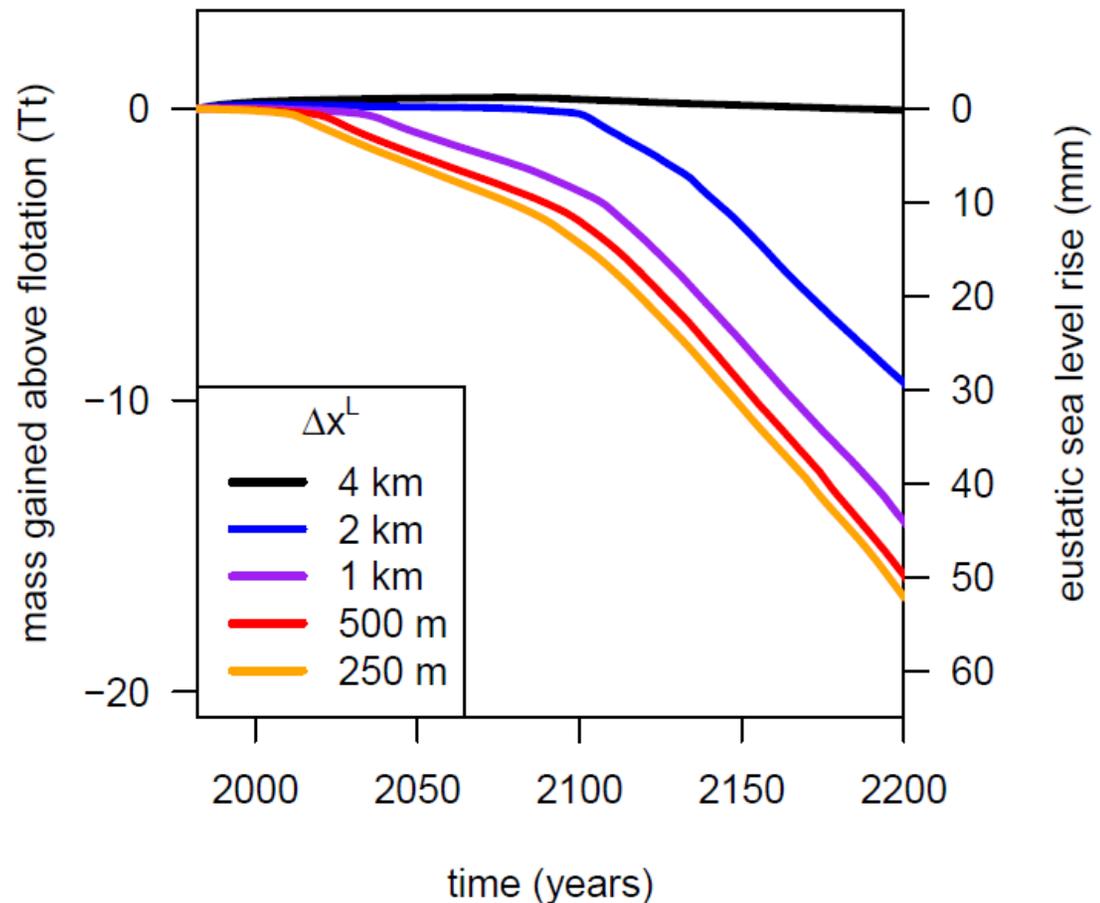
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# Ice2Sea Amundsen (cont)

- Need at least 2 km resolution to get any measurable contribution to SLR.
- Appears to converge at first-order in  $\Delta x$

SLR vs. year, Amundsen Sea Sector



# *Ice2Sea Amundsen (cont) - Thwaites?*

- In 400 year run, Thwaites destabilizes as well.
- Essentially same forcing as previous run, subshelf melting held constant past 2200.
  - More melting: extra 5 m/a
- Thwaites is very stable, *until it tips*.



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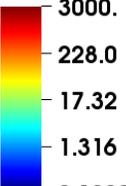
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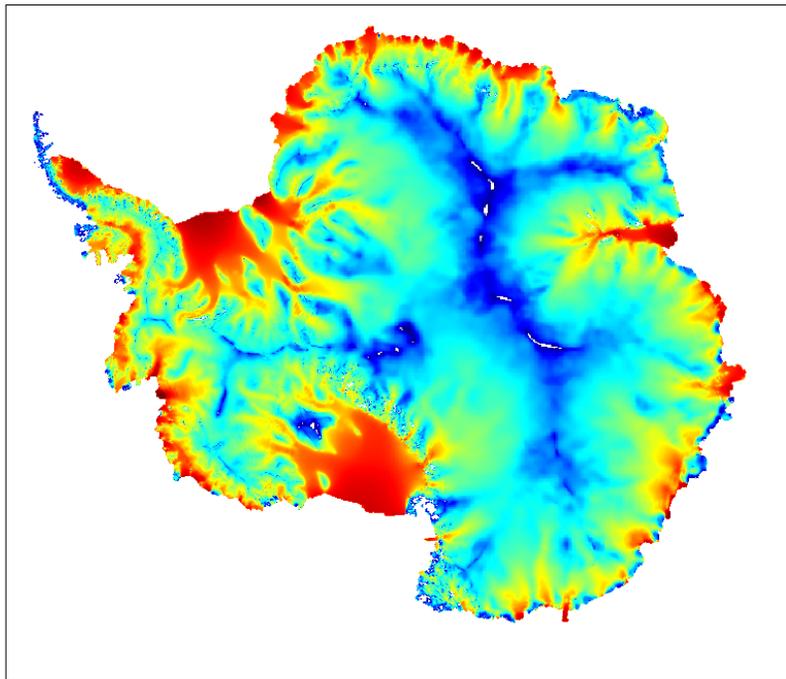
# Antarctica (Ice2Sea)

- Refinement based on Laplacian(velocity), grounding lines
- 5 km base mesh with 3 levels of refinement
  - base level (5 km): 409,600 cells (100% of domain)
  - level 1 (2.5 km): 370,112 cells (22.5% of domain)
  - Level 2 (1.25 km): 955,072 cells (14.6% of domain)
  - Level 3 (625 m): 2,065,536 cells (7.88% of domain)

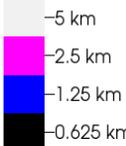
Mag(Velocity)



3000.  
228.0  
17.32  
1.316  
0.1000  
Max: 5381.  
Min: 0.000



Mesh Resolution



5 km  
2.5 km  
1.25 km  
0.625 km



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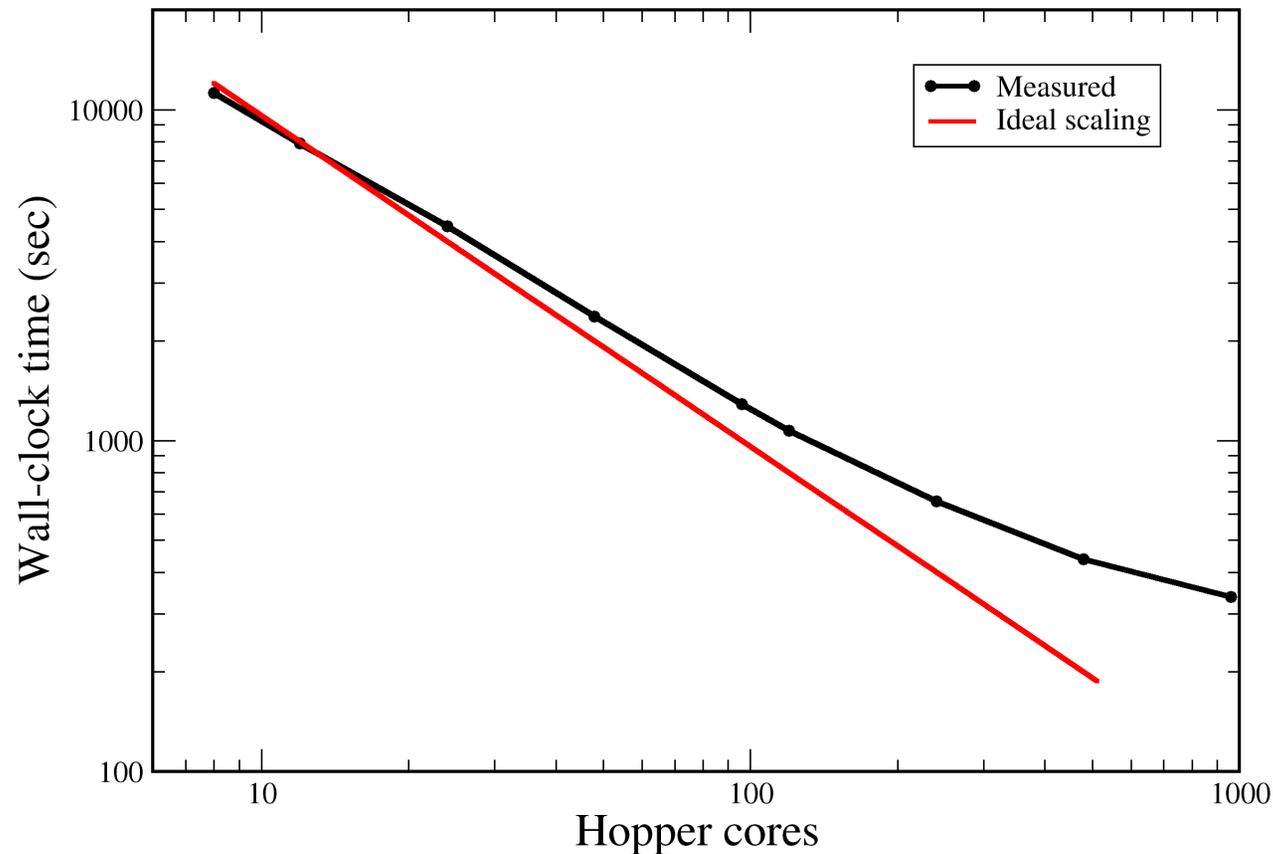


University of  
**BRISTOL**

# Parallel scaling, Antarctica benchmark

Strong Scaling of Antarctica Test Problem

[hopper.nersc.gov](http://hopper.nersc.gov)



(Preliminary scaling result – includes I/O and serialized initialization)



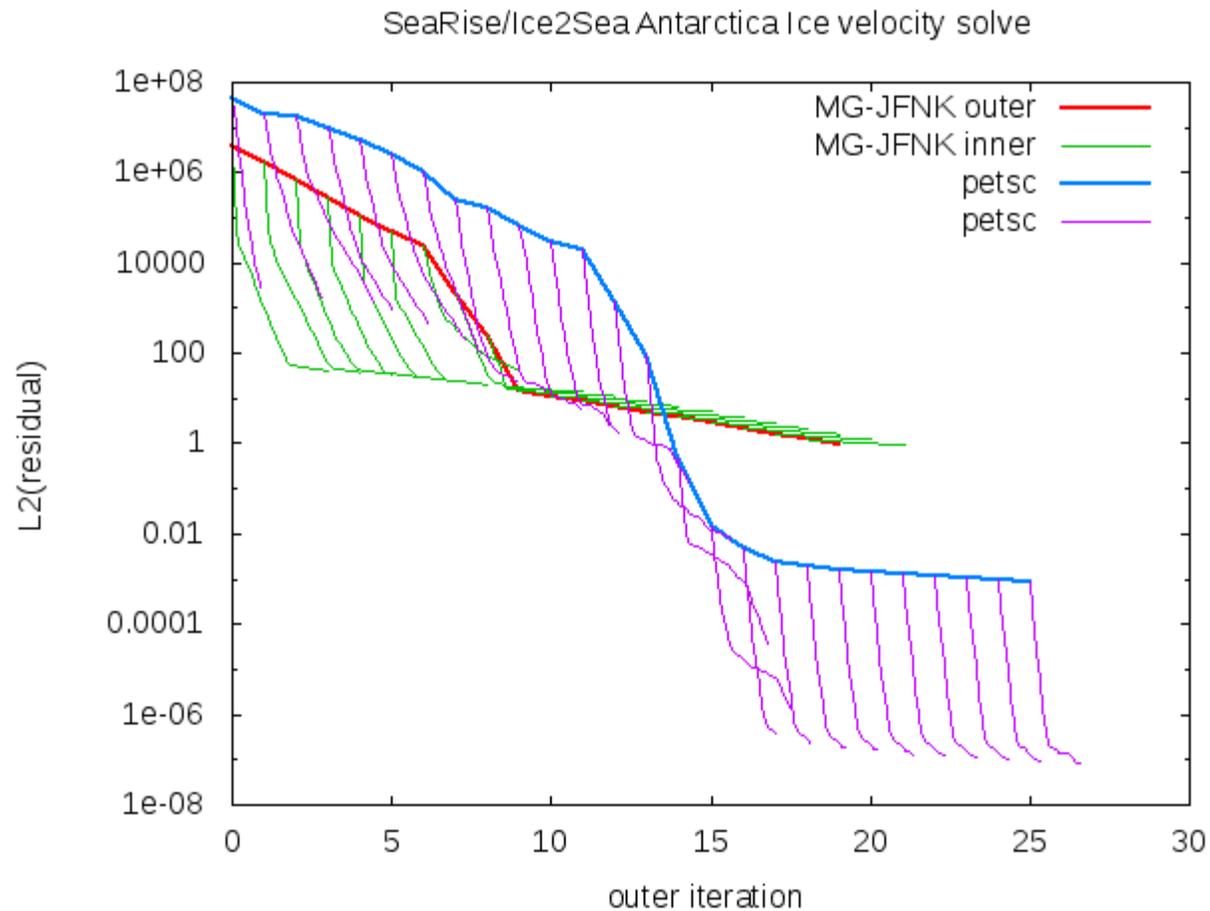
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# Linear Solvers - GAMG vs. Geometric MG



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# Temporal Discretization

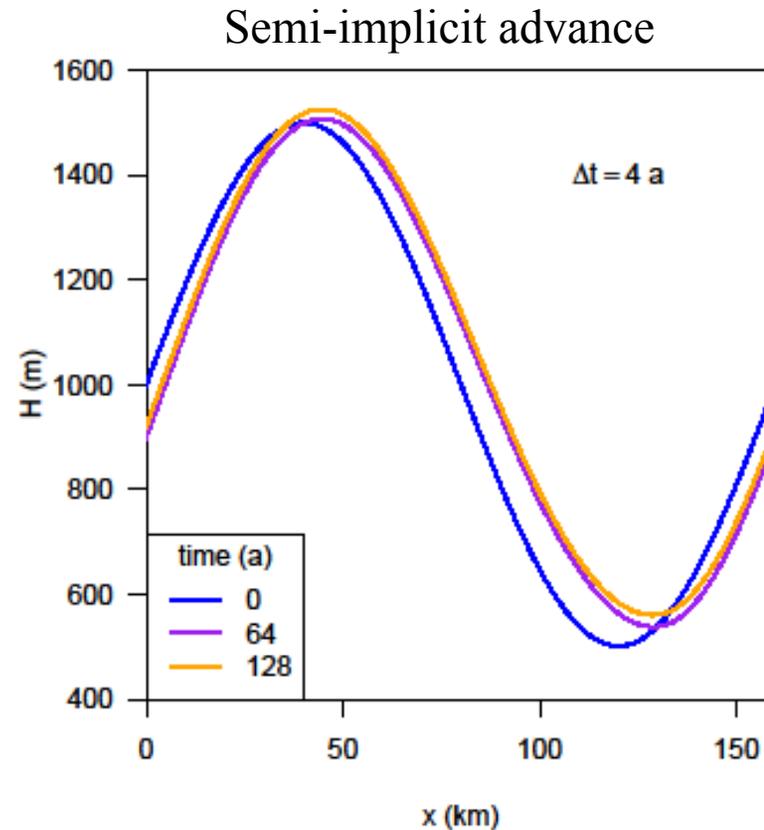
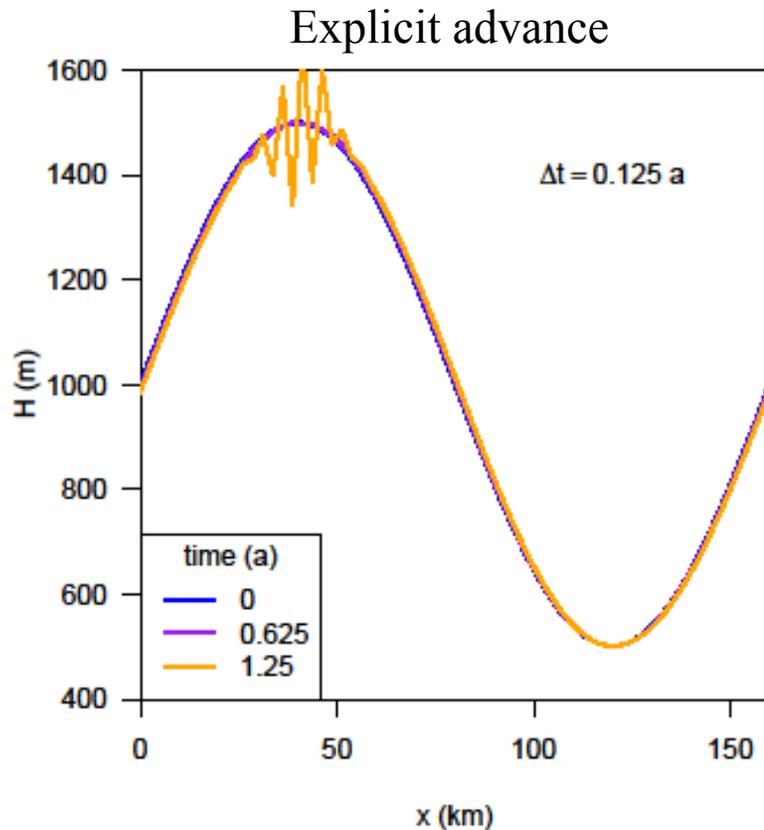
Update equation for H:  $\partial H / \partial t + \nabla \cdot (u H) = S$

- ❑ “looks” like hyperbolic advection equation (explicit scheme, Courant stability --  $\Delta t \propto \Delta x$ )
- ❑ Velocity field has  $\nabla H$  piece - diffusion equation for H ( $\Delta t \propto \Delta x^2$  !)
- ❑ Strategy (Cornford) - try to factor out diffusive flux and discretize as an advection-diffusion equation:
  - ❑  $F = u H = F \downarrow \text{advective} + F \downarrow \text{diffusive}$
  - ❑  $F \downarrow \text{diffusive} = -D \nabla H$
  - ❑ Now solve:  $\partial H / \partial t + \nabla \cdot F \downarrow \text{advective} = \nabla \cdot (D \nabla H) + S$ 
    - ❑ Advective fluxes: explicit update using unsplit 2<sup>nd</sup> Order PPM scheme
    - ❑ Diffusive fluxes: implicit update (Backward Euler for now)



# Temporal Discretization (cont)

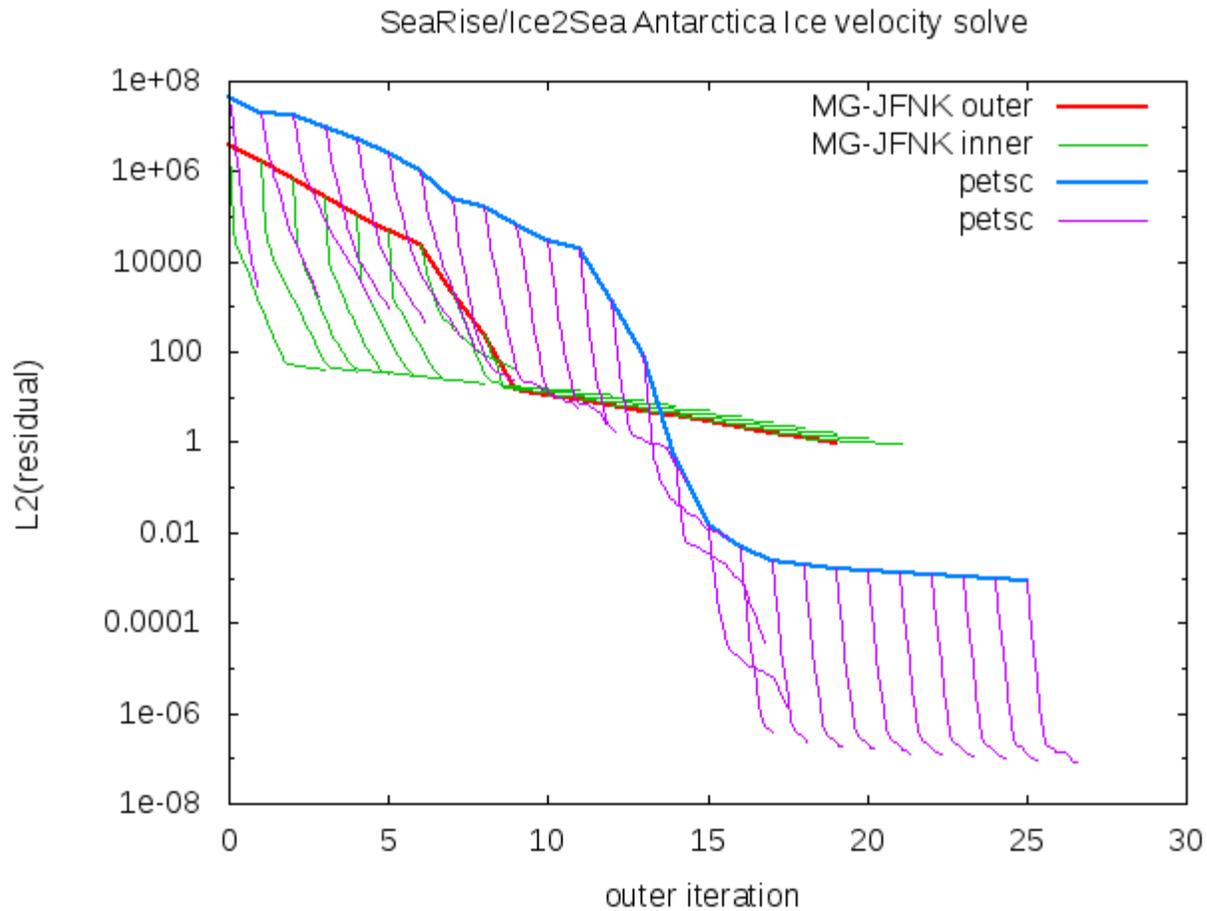
- ❑ Test case based on ISMIP-HOM A geometry
- ❑  $\Delta x = 2.5 \text{ km}$ ,  $\Delta t \downarrow CFL = 5 \text{ a}$



- ❑ Unfortunately, still run into stability issues finer than  $\Delta x < 0.5 \text{ km}$



# Linear Solvers - GAMG vs. Geometric MG



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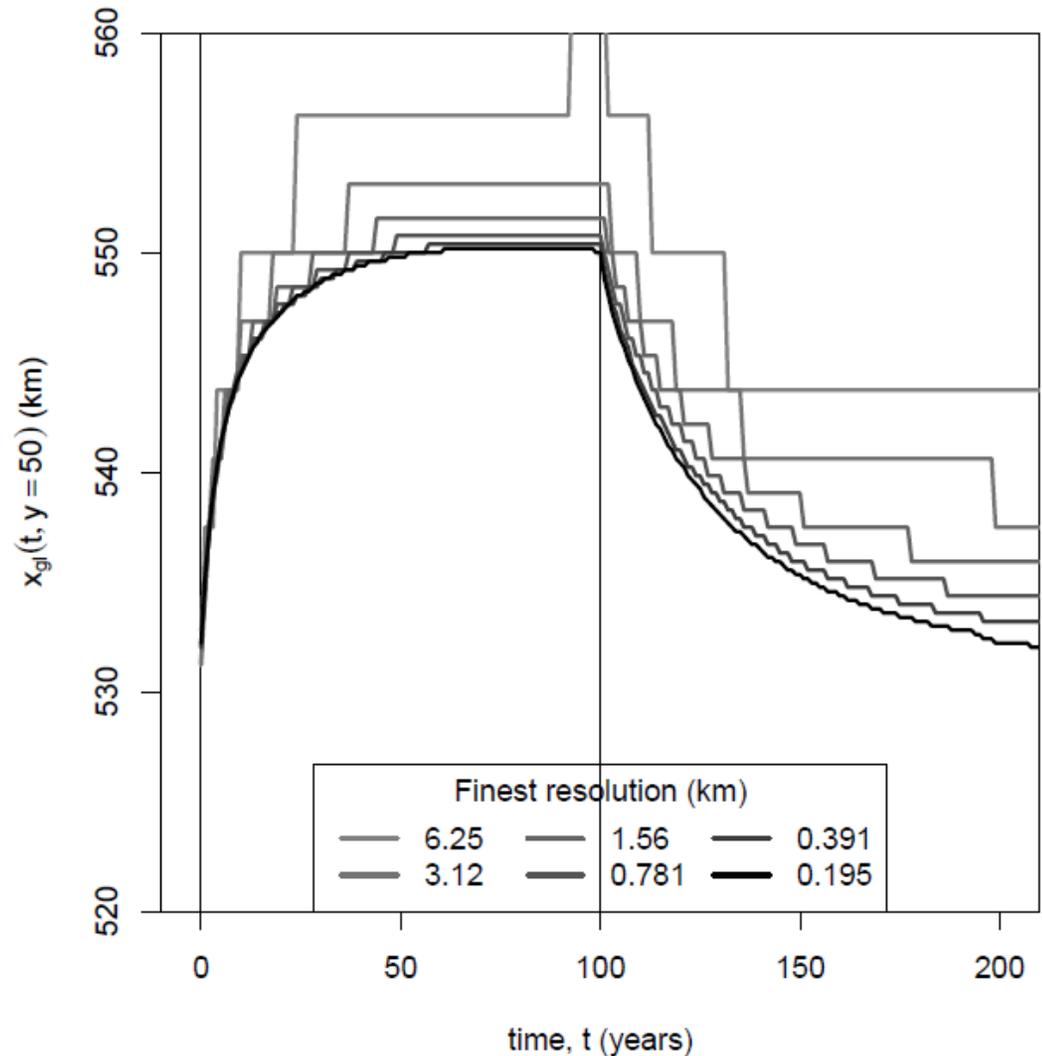
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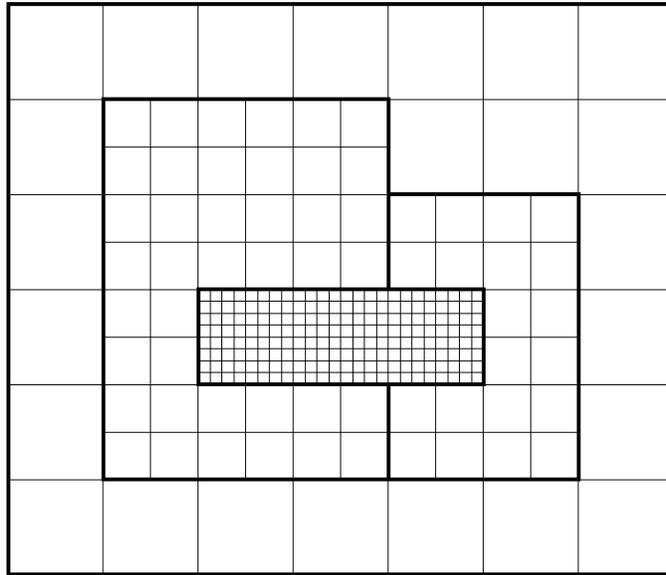
# MISMIP3D: Mesh resolution

- Plot shows grounding line position  $x \downarrow GL$  at  $y=50\text{ km}$  vs. time for different spatial resolutions.
- $\Delta x = 0.195\text{ km} \rightarrow 6.25\text{ km}$
- Appears to require finer than 1 km mesh to resolve dynamics
- Converges as  $O(\Delta x)$  (as expected)

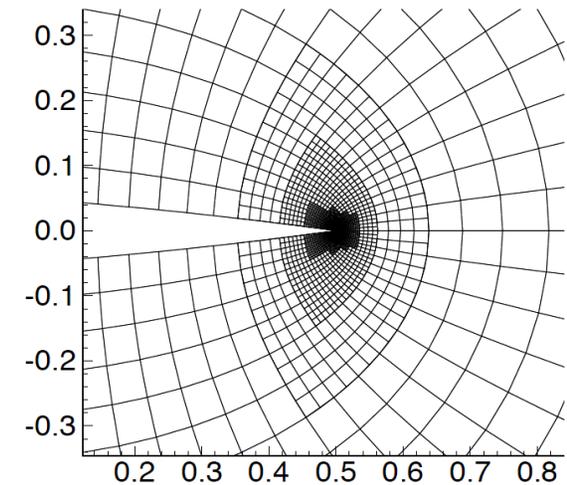
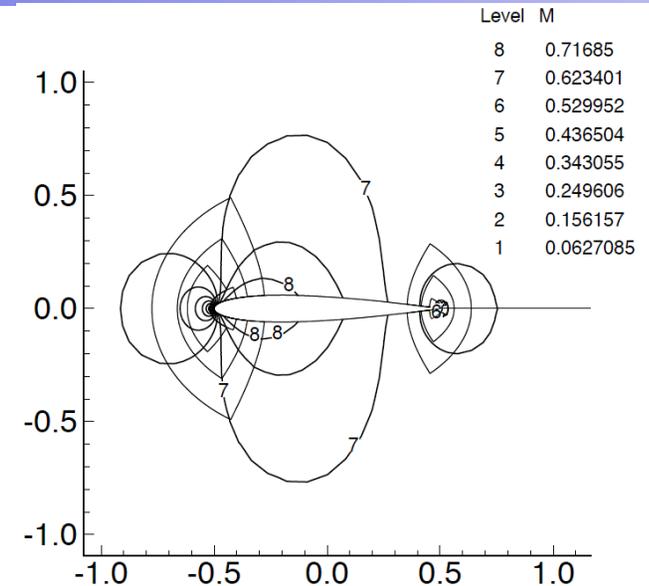


# Block-Structured Local Refinement

- Refined regions are organized into rectangular patches.

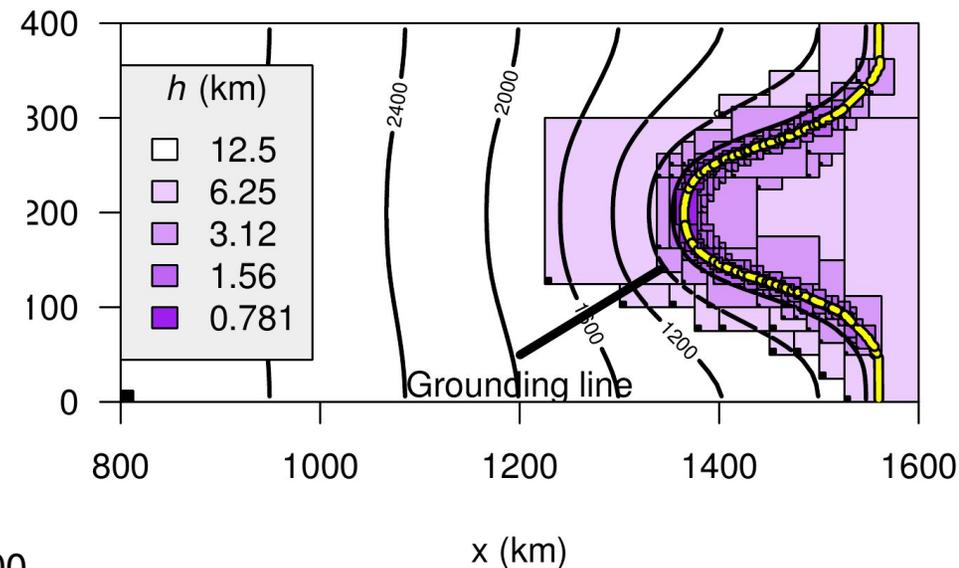
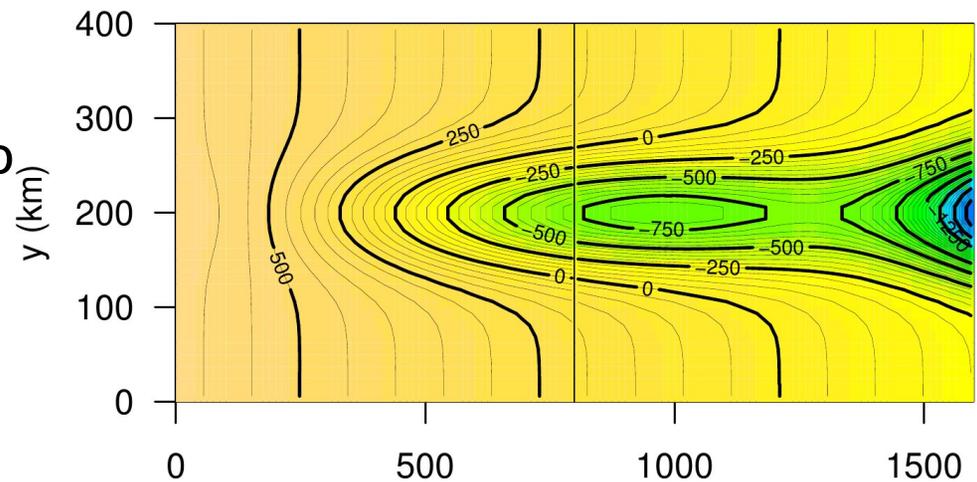
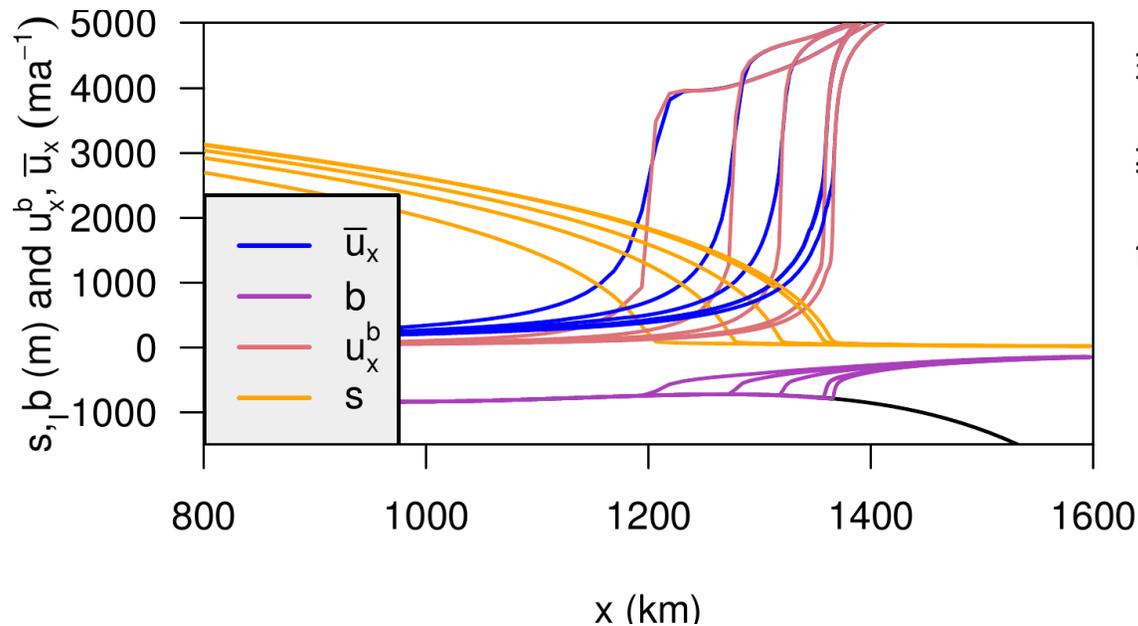


- *Algorithmic advantages:*
  - *Build on mature structured-grid discretization methods.*
  - *Low overhead due to irregular data structures, relative to single structured-grid algorithm.*



# BISICLES results - Grounding line study

- ❑ Bedrock topography based on Katz and Worster (2010)
- ❑ Evolve initially uniform-thickness ice to steady state
- ❑ Repeatedly add refinement and evolve to steady state
- ❑ G.L. advances with finer resolution
- ❑ Appear to need better than 1 km

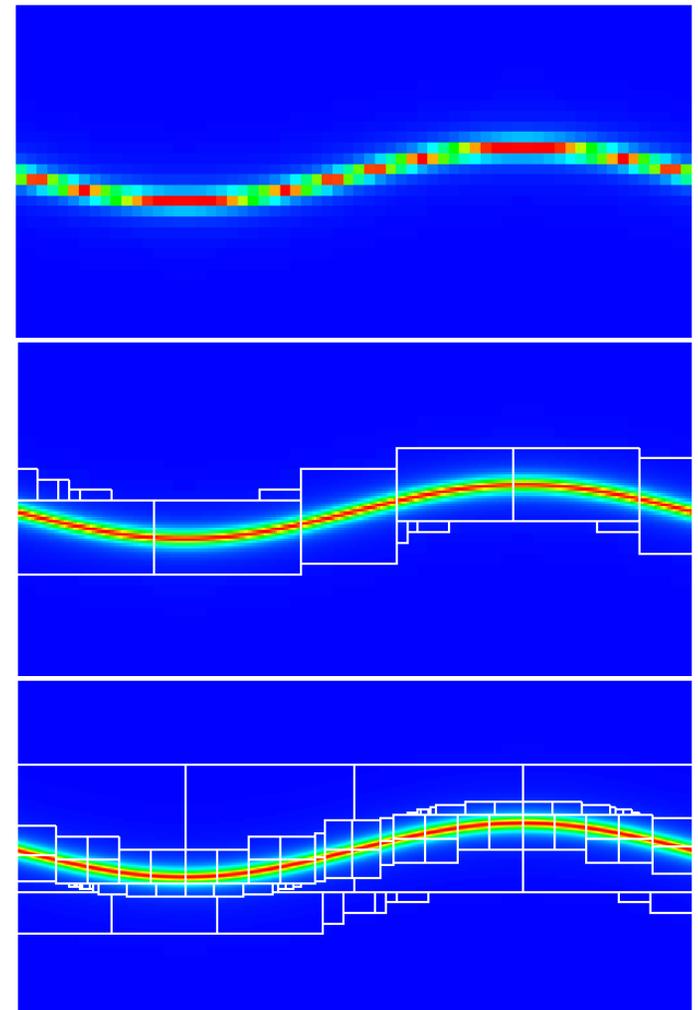
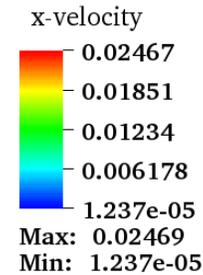


# BISICLES Results

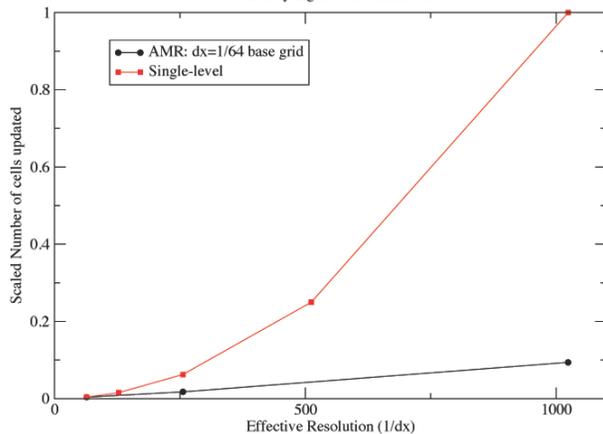
## ❑ Ice-stream Simulation

[based on Pattyn et al (2008)]:

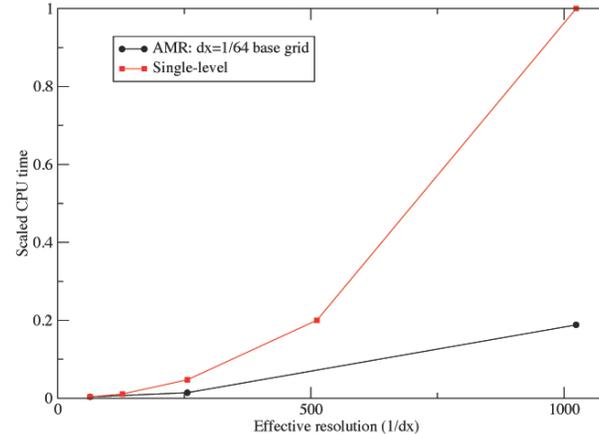
- High resolution is required to accurately resolve the ice stream.
- AMR simulation allows high resolution around the ice stream at a fraction of the cost of a uniformly refined mesh.



Number of cells updated  
Scaled by highest-resolution run



CPU Times for AMR vs. non-AMR  
Scaled by highest-resolution run



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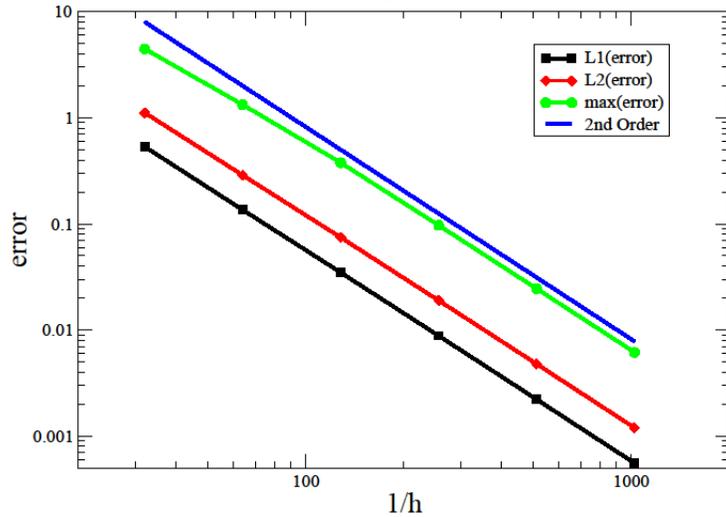
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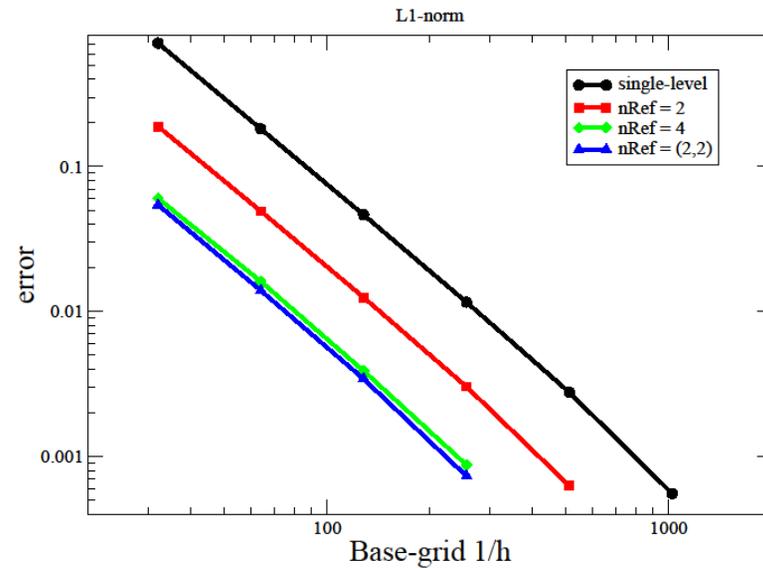


# Numerical Accuracy and Convergence

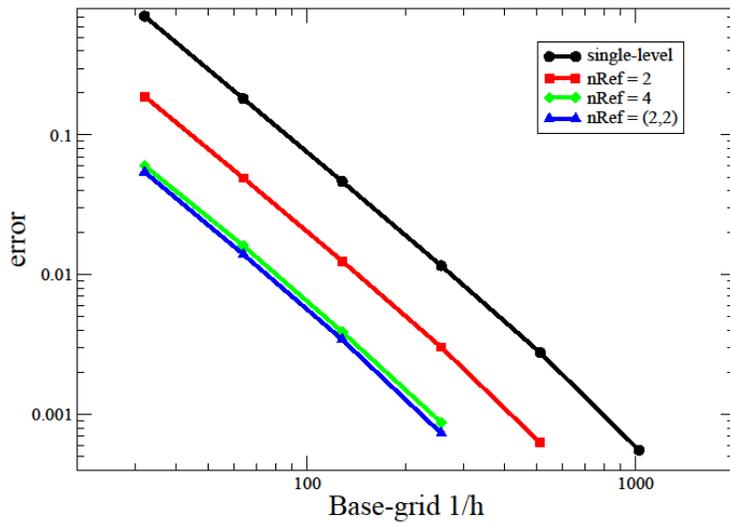
Richardson Convergence of x-velocity  
(single-level)



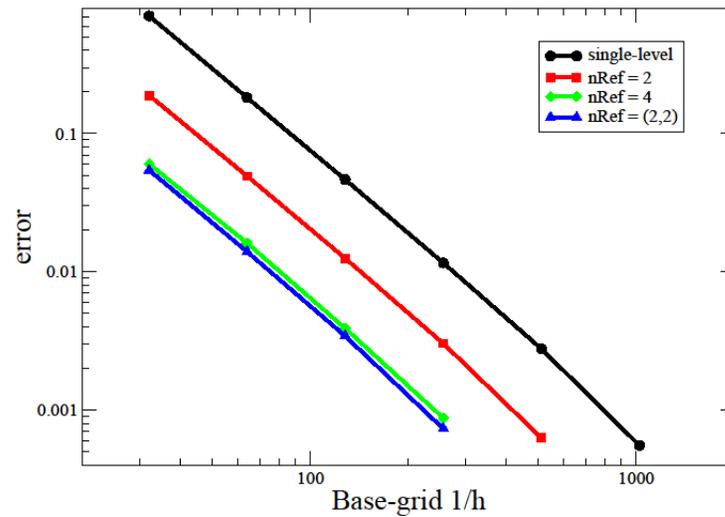
x-velocity AMR Convergence



x-velocity AMR Convergence  
L1-norm



x-velocity AMR Convergence  
L1-norm



# MISMIP+

- ❑ “Child of MISMIP3D”
  - Examined GL response of models to a localized change in bed friction
  - Clarified resolution requirements for reversible GL dynamics
  - **Large variation in steady-state GL position among models**
  - Conclusions about dynamical results clouded by this difference
  - Said nothing about response to subshelf melt forcing (buttressing?)
  
- ❑ Specific details still under development
  - Steady-state with reduced variation between models
    - Steady-state on upward-sloping bed (buttressing) -- Gudmundsson (2012)
    - Narrow-ish channel (still under discussion)
  - Perturbation due to subshelf melt anomaly - GL retreat
  - Reversibility? (return timescale seems long)
  - Primary contact - Steph Cornford (Bristol)

